# https://doi.org/10.46813/2021-134-070 PLASMA LENS FOR ELECTRON AND POSITRON BEAMS

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Focusing of both electron and positron bunches in electron-positron collider is necessary. When long electron/positron bunch is injected into the plasma, the focusing force is not uniform but oscillated. It is shown that a long positron bunch after focusing is destroyed faster than an electron bunch due to betatron and plasma oscillations. PACS: 29.17.+W; 41.75.LX

## **INTRODUCTION**

Plasma wakefield accelerators provide an extremely high accelerating gradient [1 - 13], long sequence focusing and large transformer ratio obtaining [14 - 23], resonant wakefield excitation by a nonresonant sequence of electron bunches [24, 25]. Focusing of relativistic electron bunches by wakefield, excited in the plasma, is important previously studied effect [26 - 37]. Mechanism of focusing in the plasma, by which all bunches of a sequence are focused identically and uniformly, is proposed and investigated by numerical simulation in [14 - 16].

The plasma lens for long relativistic electron and positron bunch focusing is studied in this paper by numerical simulation using code lcode [38]. Unbounded non-magnetized homogeneous plasma is considered. The Gaussian  $(\cos^2 A)$  in longitudinal direction bunch (current profile) is considered.

The purpose of this paper is to show by numerical simulation that one can achieve conditions of focusing of long relativistic electron and positron bunch. In this paper the authors present the results of numerical simulation on the focusing force distribution for long electron and positron bunches.

We use the cylindrical coordinate system (r, z). The time  $\tau$  is normalized to  $\omega_{_{pe}}^{^{-1}}$ , all the distances and coordinate  $\xi$  – to  $c\omega_{pe}^{-1}$ , the density – to the unperturbed plasma electron density, the beam current  $I_b$  – to  $mc^3 / e = 17 \text{ kA}$ , the fields – to  $mc\omega_{pe} / e$ , where m is the electron mass, e is the electron charge, c is the speed of light,  $\omega_{pe}$  is the plasma electron frequency.  $\lambda$  is the plasma wavelength. The simulation time is  $60.1\omega_{pe}^{-1}$ .  $\gamma_b = 5$  for all bunches. The length of all bunches is  $L_b = 8\lambda$ . These normalisations are used also in the Figures. The arrow on all Figures indicates the direction of movement of the bunches.

We present the numerical simulation data on plasma wakefield excitation by a relativistic electron and positron bunch, obtained with the 2.5D quasi-static code lcode. It treats the plasma as a cold electron fluid (magnetohydrodynamics model), and the bunches as ensembles of macro-particles.

Spatial step equals  $0.1c\omega_{pe}^{-1}$ . Time step for plasma electrons equals  $0.1\omega_{pe}^{-1}$ . Time step for beam electrons equals  $0.1\sqrt{\gamma_b}\omega_{pe}^{-1}$ . Spatial dependences in selected points of observation are presented.

### **RESULTS OF NUMERICAL SIMULATION**

At first, the excited field distribution, formed by long Gaussian electron bunch in the plasma (Fig. 1) is considered.



Fig. 1. Spatial distribution of Gaussian bunch electron density  $n_b(\xi, r)$  (minus indicates to the electron

bunches) at  $\gamma_b = 5$ ,  $2\sigma_z = 25$ ,  $\sigma_r = 0.1$ , maximum initial electron bunch current  $I_{b} = 0.6 \cdot 10^{-3}$ 

In this case, a rather smooth electrons  $n_e(\xi, r)$ (Fig. 2) density pit is formed in the plasma in the Gaussian bunch region.



Fig. 2. Spatial distribution of plasma electron density  $n_{e}(\xi,r)$  (corresponds to Fig. 1)

In addition, smooth focusing force is observed in the region of the bunch (Fig. 3).

First (approximately, during the first half of the simulation time), the focusing of the electron bunch is observed (Figs. 4-6). The centers of the bunches are subjected to the strongest focusing (see Fig. 5).





Fig. 4. Spatial distribution of focused Gaussian bunch electron density  $n_b(\xi, r)$ 



Fig. 5. Longitudinal distribution of bunch electron density  $n_b(\xi, r=r_b)$ : before focusing (a); after focusing (b)



Fig. 6. Longitudinal distribution of bunch electron density  $n_b(\xi, r=0)$ : before focusing (a); after focusing (b)

Then (after approximately  $t = 30\omega_{pe}^{-1} - 40\omega_{pe}^{-1}$ ) betatron oscillations develop. Since the frequency of betatron oscillations for electrons of the bunch located at different radii is different, the bunch stratifies and the electrons of the bunch peripheral along the radius are defocused (Fig. 7). In general, the destruction of the bunch can be observed along the entire diameter.



From comparison Figs. 7 and 8 it can be seen that for the same time the positron bunch is destroyed more strongly than electron.

From the beginning, oscillations are observed in the second part (after  $\xi = 30 c / \omega_{pe}$ ) of the plasma electron density pit (Fig. 9).





Then oscillations are excited in the plasma electron density (Fig. 10).

The development of oscillations can be seen on  $F_r(\xi)$  graph (Fig. 11).

As a result, the bunch is modulated (see Fig. 12).



Fig. 12. Longitudinal distribution of bunch electron density  $n_b(\xi, r = r_b)$  at large (end of simulation) times

In general, the destruction of the bunch is visible along the entire diameter. Thus, the inhomogeneity of the focusing force and the modulation of relativistic bunches of electrons and positrons during their propagation in the plasma due to betatron and plasma oscillations are shown. In addition, it has been shown that positron bunch is destroyed more strongly and faster than electron bunch.

#### CONCLUSIONS

A numerical simulation of the focusing of electron and positron bunches by a plasma lens is carried out. When a long electron/positron bunch is injected into the plasma, the resulting focusing force is not uniform, but oscillates. It was shown that a long bunch of positrons after focusing is destroyed faster than an electron bunch due to betatron and plasma oscillations.

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# ПЛАЗМЕННАЯ ЛИНЗА ДЛЯ ЭЛЕКТРОННОГО И ПОЗИТРОННОГО ПУЧКОВ

## Д.С. Бондарь, В.И. Маслов, И.Н. Онищенко, Р.Т. Овсянников

Необходима фокусировка как электронных, так и позитронных сгустков в электрон-позитронных коллайдерах. При инжекции длинного сгустка электронов/позитронов в плазму образующаяся фокусирующая сила не однородна, а с некоторыми осцилляциями. Показано, что длинный позитронный сгусток после фокусировки разрушается быстрее, чем электронный сгусток за счет бетатронных колебаний и плазменных осцилляций.

# ПЛАЗМОВА ЛІНЗА ДЛЯ ЕЛЕКТРОННОГО І ПОЗИТРОННОГО ПУЧКІВ

# Д.С. Бондарь, В.І. Маслов, І.М. Оніщенко, Р.Т. Овсянніков

Необхідне фокусування як електронних, так і позитронних згустків в електрон-позитронних колайдерах. При інжекції довгого згустка електронів/позитронів у плазму утворювана фокусуюча сила не однорідна, а з деякими осциляціями. Показано, що довгий позитронний згусток після фокусування руйнується швидше, ніж електронний згусток за рахунок бетатронних коливань і плазмових осциляцій.