

INTEGRAL REPRESENTATION OF HYPERPARABOLIC EQUATION

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In this work, for functions that satisfy a Cauchy problem for hyperparabolic equations, we write integral equations and solve them using the method of successive approximations.

AMS Subject Classification: 60J30, 35Q40, 45L10

Introduction

The Riemman method is a well-known method for solving a Cauchy problem for the telegraph equation [1]. But this method can't be used for an analogue of the telegraph equation in R^n (hyperparabolic equations), and also nonlinear hyperparabolic equations that are defined later.

1. Cauchy Problem for Telegraph Equation

The telegraph process and the corresponding random evolution is a model for the motion of a physical particle on a line, when the particle changes its direction of motion to the opposite in random moments of time. Besides, the telegraph process defines a "rectangular wave" as an oscillating process in problems of radioengineering [2].

It's well-known [3] that a special form of a functional of the Markov random evolution satisfies the Cauchy problem

$$\frac{\partial^2}{\partial t^2} u(x, t) = -2\lambda \frac{\partial}{\partial t} u(x, t) + V^2 \frac{\partial^2}{\partial x^2} u(x, t), \quad (1)$$

$$u(x, 0) = f(x), \quad \left. \frac{\partial}{\partial t} u(x, t) \right|_{t=0} = V \frac{d}{dx} f(x).$$

But, for the functional $u(x, t)$, it is possible to write an integral equation from the first jump of the process,

$$u(x, t) = e^{-\lambda t} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t f(x - V(t + 2s)) ds + \lambda^2 \int_0^t \int_0^{t-s} e^{-\lambda(l+s)} u(x + V(s-l), t - (l+s)) dl ds. \quad (2)$$

The problem (1) can be obtained from (2). To do this, let us change the variables in (2),

$$u(x, t) = e^{-\lambda t} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t f(x - V(t + 2s)) ds$$

$$- \lambda^2 e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} u(x + V(2s + f - t), f) dsdf.$$

Differentiating the last expression with respect to x and t we have

$$\frac{\partial^2}{\partial x^2} u(x, t) = e^{-\lambda t} \frac{d^2}{dx^2} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t \frac{d^2}{dx^2} f(x - V(t + 2s)) ds$$

$$- \lambda^2 e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} \frac{\partial^2}{\partial x^2} u(x + V(2s + f - t), f) dsdf,$$

$$\frac{\partial}{\partial t} u(x, t) = -\lambda \left[e^{-\lambda t} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t f(x - V(t - 2s)) ds \right.$$

$$\left. - \lambda^2 e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} u(x + V(2s + f - t), f) dsdf \right] + \left[V e^{-\lambda t} \frac{d}{dx} f(x + Vt) \right.$$

$$\left. - \lambda V e^{-\lambda t} \int_0^t \frac{d}{dx} f(x - V(t - 2s)) ds + \lambda^2 V e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} \right.$$

$$\left. \times \frac{\partial}{\partial x} u(x + V(2s + f - t), f) dsdf \right] + \lambda e^{-\lambda t} f(x + Vt) - \lambda^2 e^{-\lambda t} \int_0^t e^{\lambda f}$$

$$\times u(x + Vt - Vf, f) df = -\lambda u(x, t) + \left[V e^{-\lambda t} \frac{d}{dx} f(x + Vt) \right.$$

$$\left. - \lambda V e^{-\lambda t} \int_0^t \frac{d}{dx} f(x - V(t - 2s)) ds + \lambda^2 V e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} \right.$$

$$\begin{aligned}
& \times \left. \frac{\partial}{\partial x} u(x + V(2s + f - t), f) dsdf \right] + \lambda e^{-\lambda t} f(x + Vt) \\
& - \lambda^2 e^{-\lambda t} \int_0^t e^{\lambda f} u(x + Vt - Vf, f) df, \\
\frac{\partial^2}{\partial t^2} u(x, t) = & - \lambda \frac{\partial}{\partial t} u(x, t) - \lambda \left[V e^{-\lambda t} \frac{d}{dx} f(x + Vt) - \lambda V e^{-\lambda t} \right. \\
& \times \int_0^t \frac{d}{dx} f(x - V(t - 2s)) ds + \lambda^2 V e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} \frac{\partial}{\partial x} u(x + V(2s \\
& + f - t), f) dsdf + \lambda e^{-\lambda t} f(x + Vt) - \lambda^2 e^{-\lambda t} \int_0^t e^{\lambda f} u(x + Vt - Vf, f) df \left. \right] \\
& + V^2 \left[e^{-\lambda t} \frac{d^2}{dx^2} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t \frac{d^2}{dx^2} f(x - V(t - 2s)) ds \right. \\
& \left. - \lambda^2 e^{-\lambda t} \int_0^t \int_0^{t-f} e^{\lambda f} \frac{\partial^2}{\partial x^2} u(x + V(2s + f - t), f) dsdf \right] - \lambda V e^{-\lambda t} \frac{d}{dx} f(x \\
& + Vt) + \lambda^2 V e^{-\lambda t} \int_0^t e^{\lambda f} \frac{d}{dx} u(x + Vt - Vf, f) df + \lambda V e^{-\lambda t} \frac{d}{dx} f(x \\
& + Vt) - \lambda^2 V e^{-\lambda t} \int_0^t e^{\lambda f} \frac{d}{dx} u(x + Vt - Vf, f) df + \lambda^2 e^{-\lambda t} e^{\lambda t} u(x, t) \\
= & - \lambda \frac{\partial}{\partial t} u(x, t) - \lambda \left[\frac{\partial}{\partial t} u(x, t) + \lambda u(x, t) \right] + V^2 \frac{\partial^2}{\partial x^2} u(x, t) + \lambda^2 u(x, t) \\
= & - 2\lambda \frac{\partial}{\partial t} u(x, t) + V^2 \frac{\partial^2}{\partial x^2} u(x, t).
\end{aligned}$$

We thus obtained the Cauchy problem (1). The following theorem holds.

Theorem 1. *The function $u(x, t)$, satisfying (1) under the condition of integrability of $u(x, t)$ and $f(x)$, satisfies equation (2). The function $u(x, t)$, satisfying (1) under the condition of differentiability with respect to x and t , satisfies the Cauchy problem (1).*

The Cauchy problem (1) and equation (2) are equivalent in this sense.

Let us consider the question of whether there exist a solution of (2) in the space of functions

$$\phi(x, t) = \phi_0(x, t) + c, \tag{3}$$

where $c = \text{const}$, $\phi_0(x) \rightarrow 0, x, t \rightarrow \infty$. This Banach space with sup-norm was studied in the works of V.S.Koroljuk and A.F.Turbin (see, for example, [4]).

Let us write (2) in the following form: $u(x, t) = Au(x, t)$, where $Au(x, t)$ is equal the right-hand side of (2). A acts in the space (3), when $f(x) = f_0(x) + c$. Indeed,

$$\begin{aligned} A(\phi_0 + c) &= e^{-\lambda t}(f_0(x) + c) + \lambda e^{-\lambda t} \int_0^t [f_0(x)(x - V(t - 2s)) + c] ds \\ &+ \lambda^2 \int_0^t \int_0^{t-s} e^{-\lambda(l+s)} [\phi_0(x + V(s - l), t - (s + l)) + c] dlds = \left\{ f_0(x)e^{-\lambda t} \right. \\ &+ \lambda e^{-\lambda t} \int_0^t f_0(x)(x - V(t - 2s)) ds + \lambda^2 \int_0^t \int_0^{t-s} e^{-\lambda(l+s)} \phi_0(x + V(s - l), t \\ &- (s + l)) dlds \left. \right\} + \left\{ e^{-\lambda t} c + \lambda e^{-\lambda t} t c + \lambda^2 \left(-\frac{1}{\lambda} e^{-\lambda t} c t - \frac{1}{\lambda^2} (e^{-\lambda t} - 1) c \right) \right\} \\ &= \left\{ f_0(x)e^{-\lambda t} + \lambda e^{-\lambda t} \int_0^t f_0(x)(x - V(t - 2s)) ds + \lambda^2 \int_0^t \int_0^{t-s} e^{-\lambda(l+s)} \right. \\ &\quad \left. \times \phi_0(x + V(s - l), t - (s + l)) dlds \right\} + c, \end{aligned}$$

where the expression in braces converges to 0 as $x, t \rightarrow \infty$. Indeed,

$$e^{-\lambda t} \int_0^t f(x) dx < e^{-\lambda t} \int_0^t \sup_x |f(x)| dx = \sup_x |f(x)| t e^{-\lambda t} \rightarrow 0,$$

$$x, t \rightarrow \infty, t > 0,$$

where k is even in the second sum,

$$\lim_{n \rightarrow \infty} u_n(x, t) = x + \frac{V}{2\lambda}(1 - e^{-2\lambda t}),$$

which coincides with the first moment of the Markov random evolution found in [5].

It should be noted that this method can be used for hyperparabolic equations in R^n .

2. Nonlinear Hyperparabolic Equation

A fading telegraph process and a fading Markov random evolution defines a motion of a particle on a line in the field of gravity, if the particle is attracted to some point on the line. From the point of view of radioengineering a fading telegraph process defines a fading “rectangular wave”.

When we consider a generalization of a telegraph process, a fading telegraph process, another integral equation appears,

$$\begin{aligned} u(V, x, t) = & e^{-\lambda t} f(x + Vt) + \lambda e^{-\lambda t} \int_0^t f\left(x + Vs - \frac{V}{c}(t - s)\right) \\ & + \lambda^2 \int_0^t \int_0^{t-s} e^{-\lambda(l+s)} u\left(\frac{V}{c^2}, x + Vs - \frac{V}{c}l, t - (s + l)\right) dl ds. \end{aligned} \quad (4)$$

Let us note that for $c = 1$, equation (4) coincides with (2), where $u(V, x, t) = u(x, t)$.

By making the changes of variables, similar to the one in Section 1, and differentiating we get a Cauchy problem corresponding to (4),

$$\begin{aligned} & \frac{\partial^3}{\partial t^3} u(V, x, t) + 3\lambda \frac{\partial^2}{\partial t^2} u(V, x, t) + 3\lambda^2 \frac{\partial}{\partial t} u(V, x, t) - \lambda^2 \frac{\partial}{\partial t} u\left(\frac{V}{c^2}, x, t\right) \\ & - \frac{V^2}{c^2} \frac{\partial^3}{\partial x^2 \partial t} u(V, x, t) - \frac{V^2 \lambda}{c^2} \frac{\partial^2}{\partial x^2} u(V, x, t) \\ & - \frac{c-1}{c} \lambda^2 V u\left(\frac{V^2}{c^2}, x, t\right) + \lambda^3 u(V, x, t) - \lambda^3 u\left(\frac{V}{c^2}, x, t\right) = 0, \end{aligned} \quad (5)$$

$$u(V, x, t) = f(x), \quad \frac{\partial}{\partial t} u(x, t) \Big|_{t=0} = V \frac{d}{dx} f(x),$$

$$\frac{\partial^2}{\partial t^2} u(x, t) \Big|_{t=0} = -\frac{c+1}{c} \lambda V \frac{d}{dx} f(x) + V^2 \frac{d^2}{dx^2} f(x).$$

For $c = 1$, we have

$$\begin{aligned} & \frac{\partial^3}{\partial t^3} u(x, t) + 3\lambda \frac{\partial^2}{\partial t^2} u(x, t) + 2\lambda^2 \frac{\partial}{\partial t} u(x, t) \\ & - V^2 \frac{\partial^3}{\partial x^2 \partial t} u(x, t) - V^2 \lambda \frac{\partial^2}{\partial x^2} u(x, t) = 0 \end{aligned} \quad (6)$$

or

$$\begin{aligned} & \left\{ \frac{\partial^3}{\partial t^3} u(x, t) + 2\lambda \frac{\partial^2}{\partial t^2} u(x, t) - V^2 \frac{\partial^3}{\partial x^2 \partial t} u(x, t) \right\} \\ & + \left\{ \lambda \frac{\partial^2}{\partial t^2} u(x, t) + 2\lambda^2 \frac{\partial}{\partial t} u(x, t) - V^2 \lambda \frac{\partial^2}{\partial x^2} u(x, t) \right\} = 0, \\ & \frac{\partial}{\partial t} \left\{ \frac{\partial^2}{\partial t^2} u(x, t) + 2\lambda \frac{\partial}{\partial t} u(x, t) - V^2 \frac{\partial^2}{\partial x^2} u(x, t) \right\} \\ & + \lambda \left\{ \frac{\partial^2}{\partial t^2} u(x, t) + 2\lambda \frac{\partial}{\partial t} u(x, t) - V^2 \frac{\partial^2}{\partial x^2} u(x, t) \right\} = 0, \\ & \left\{ \frac{\partial}{\partial t} + \lambda \right\} \left\{ \frac{\partial^2}{\partial t^2} + 2\lambda \frac{\partial}{\partial t} - V^2 \frac{\partial^2}{\partial x^2} \right\} u(x, t) = 0. \end{aligned}$$

This is a factorized equation, one component of which coincides with (1). Correspondingly, if $u(x, t)$ satisfies (5) for $c = 1$, then it satisfies (1).

The proof of existence of a solution of (4) is similar to that in Section 1 for the space of functions $\phi(V, x, t) = \phi_0(V, x, t) + c$. The following theorem holds.

Theorem 3. *The Cauchy problem for nonlinear hyperparabolic equation (5) is equivalent to integral equation (4) that has a solution in the space of functions $\phi(V, x, t) = \phi_0(V, x, t) + c$, where $c = \text{const}$, $\phi_0(V, x, t) \rightarrow 0$, $V, x, t \rightarrow \infty$.*

As in Section 1, the method of successive approximations converges for the functions $f(x) = x^k$.

Example 2. $f(x) = x$,

$$u_0(V, x, t) = 0, u_1(V, x, t) = e^{-\lambda t} \left(x + xt + Vt - \frac{Vt^2}{2c} + \frac{Vt^2}{2} \right),$$

$$u_2(V, x, t) = u_1(V, x, t) + \lambda^2 e^{-\lambda t} \left(\frac{xt^2}{2} + \frac{xt^3}{6} + \frac{\lambda^2 t^3 V c^2 - c + 1}{6 c^2} + \frac{\lambda^3 t^4 V c^3 - c^2 + c - 1}{24 c^3} \right),$$

$$\lim_{n \rightarrow \infty} u_n(V, x, t) = x + \frac{Vc}{\lambda(c+1)} \left(1 - e^{-\lambda t \frac{1+c}{c}} \right).$$

For $c = 1$, we have $x + \frac{V}{2\lambda} (1 - e^{-2\lambda t})$ (see Example 1).

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Received 05.12.2000