Nguyen Manh Hung, Tran Thi Loan (Hanoi Pedagogical Inst., Vietnam)

## ON THE ASYMPTOTIC BEHAVIOUR OF SOLUTIONS OF THE FIRST INITIAL BOUNDARY-VALUE PROBLEMS FOR PARABOLIC EQUATIONS

## ПРО АСИМПТОТИЧНУ ПОВЕДІНКУ РОЗВ'ЯЗКІВ ПЕРШИХ ПОЧАТКОВИХ КРАЙОВИХ ЗАДАЧ ДЛЯ ПАРАБОЛІЧНИХ РІВНЯНЬ

We consider the first initial boundary-value problem for a strongly parabolic system on infinite cylinder with non-smooth boundary. We prove some results on the existence, uniqueness and asymptotic behaviour of solutions as  $t \to +\infty$ .

Розглянуто першу початкову крайову задачу для сильно параболічної системи на некінченному циліндрі із негладкою межею. Доведено деякі результати про існування, єдиність та асимптотичну поведінку розв'язків при  $t \to +\infty$ .

1. Introduction. The theory of general boundary-value problems for elliptic systems with smooth and non-smooth boundary is nearly completely considered in [1, 2]. The strongly parabolic systems on a bounded set  $\Omega_T = \Omega \times [0,T]$ ,  $\Omega \subset \mathbb{R}^n$ , with non-smooth boundary were studied in [3, 4]. However, the boundary-value problem for a parabolic system on cylinder with non-smooth boundary has been rarely considered. O. A. Ladyzhenskaya [5] showed some sufficient conditions for the existence and uniqueness of a generalized solution to the first boundary-value problem for strongly parabolic systems on a finite cylinder  $\Omega_T$ . The smoothness of generalized solutions for a parabolic equation of order 2 is considered in [6]. The general boundary-value problem on the finite cylinder for a parabolic system in the sense of Petrovski is studied in [7], where the author obtained some results on the uniqueness of a solution in the Sobolev weighted space. In [8], the second boundary-value problem for a parabolic equation of order 2 in infinite cylinder  $\Omega_\infty = \Omega \times [-\infty, +\infty)$  is considered. Some results on the asymptotic behaviour of solutions as  $t \to +\infty$  was shown in that paper.

Suppose that  $\Omega \subset \mathbb{R}^n$  is a bounded domain with its boundary  $\partial \Omega$ . Let us introduce some notation:  $\Omega_b = \Omega \times [0,b], \ \Gamma_b = \partial \Omega \times [0,b], \ \Omega_\infty = \Omega \times [0,+\infty), \ \Gamma_\infty = \partial \Omega \times [0,\infty); \ x = (x_1,\ldots,x_n) \in \Omega; \ u(x,t) = (u_1(x,t),\ldots,u_s(x,t))$  is a vector complex function;  $\alpha = (\alpha_1,\ldots,\alpha_n), \ \alpha_i \in \mathbb{N}, \ |\alpha| = \sum_{i=1}^n \alpha_i; \ D^\alpha = \partial^{|\alpha|}/\partial x_1^{\alpha_1} \ldots \partial x_n^{\alpha_n}, \ |D^\alpha u|^2 = \sum_{i=1}^s |D^\alpha u_i|^2, \ dx = dx_1 \ldots dx_n, \ u_{t^j} = (\partial^j u_1/\partial t^j,\ldots,\partial^j u_s/\partial t^j); \ \mathring{C}^\infty(\Omega)$  is the set of infinitely differentiable functions having their supports compactly embedded in  $\Omega$ ;  $H^\ell(\Omega)$  — the space consisting of functions u(x) which have generalized derivatives  $D^\alpha u_i$  belonging to  $L_2(\Omega), \ |\alpha| \leq \ell, \ 1 \leq i \leq s$  and

$$||u||_{H^{\ell}(\Omega)}^{2} = \sum_{|\alpha|=0}^{\ell} \int_{\Omega} \sum_{i=1}^{s} |D^{\alpha}u_{i}|^{2} dx < \infty;$$

 $\overset{\circ}{H}^{\ell}(\Omega)$  is the completion of  $\overset{\circ}{C}^{\infty}(\Omega)$  in the norm of  $H^{\ell}(\Omega)$ ;  $H^{\ell,k}(\Omega_T)$  is the space consisting of functions u(x,t) such that  $D^{\alpha}u_i\in L_2(\Omega_T),\ \partial^ju_i/\partial t^j\in L_2(\Omega_T),\ |\alpha|\leq \ell,\ 1\leq i\leq s,\ 1\leq j\leq k,$  with the norm

$$\|u\|_{H^{\ell,k}(\Omega_T)}^2 = \sum_{|\alpha|=0}^\ell \int\limits_{\Omega_T} |D^\alpha u|^2 dx dt + \sum_{j=1}^k \int\limits_{\Omega_T} |u_{t^j}|^2 dx dt;$$

 $H^{\ell,0}(\Omega_T)$  is the space consisting of functions u(x,t) with the norm

$$||u||_{H^{\ell,0}(\Omega_T)}^2 = \sum_{|\alpha|=0}^{\ell} \int_{\Omega_T} |D^{\alpha}u|^2 dx dt;$$

 $\overset{\circ}{H}^{\ell,k}(\Omega_T)$  is the completion of  $\overset{\circ}{C}^{\infty}(\Omega_T)$  in the norm of  $H^{\ell,k}(\Omega_T)$ ;  $\overset{\circ}{\hat{H}}^{m,0}(\Omega_T) = \{\eta(x,t) \in \overset{\circ}{H}^{m,0}(\Omega_T) : \eta(x,T) = 0\}$ . Consider the differential operator

$$L(x,t,D) = \sum_{|p|,|q|=1}^{m} D^{p} a_{pq}(x,t) D^{q} + \sum_{|p|=1}^{m} a_{p}(x,t) D^{p} + a(x,t),$$

where  $a_{pq}$ ,  $a_p$ , a are  $(s \times s)$ -matrices,  $a_{pq} = (-1)^{|p|+|q|} a_{qp}^*$ , and their elements are bounded complex functions on  $\bar{\Omega}_{\infty}$ . Moreover, for every non-zero vector  $\xi \in \mathbb{R}^n$  and  $\eta \in \mathbb{C}^s$  we have

$$\sum_{|p|,|q|=m} a_{pq}(x,t)\xi^{p}\xi^{q}\eta\bar{\eta} > 0 \quad \forall (x,t) \in \bar{\Omega}_{\infty}, \quad (1)$$

where  $\xi^p = \xi_1^{p_1} \dots \xi_n^{p_n}, \ \xi^q = \xi_1^{q_1} \dots \xi_n^{q_n}$ .

We set

$$B(u,v) = \sum_{|p|,|q|=1}^{m} (-1)^{|p|} \int_{\Omega} a_{pq} D^{q} u \overline{D^{p} v} dx$$

and

$$B_1(u,u) = B(u,u) + 2\operatorname{Re} \sum_{|p|=1}^m \int\limits_{\Omega} a_p D^p u \bar{u} dx.$$

The following lemma is proved in [9].

**Lemma 1.** If  $a_{pq}$  satisfy (1) and  $a_{pq}$  are continuous on  $\tilde{\Omega}_{\infty}$  whenever |p| = |q| = m then there exist positive numbers  $\mu_0$ ,  $\mu_1$ ,  $\lambda_0$ ,  $\lambda_1$  such that

$$(-1)^m B(u, u) \ge \mu_0 ||u||^2_{H^m(\Omega)} - \lambda_0 ||u||^2_{L^2(\Omega)},$$
 (2)

$$(-1)^m B_1(u, u) \ge \mu_1 ||u||_{H^m(\Omega)}^2 - \lambda_1 ||u||_{L^2(\Omega)}^2$$
(3)

for all  $u \in \mathring{H}^m(\Omega)$ .

The existence of generalized solution. In this section we consider the existence and uniqueness of a solution of the problem

$$(-1)^{m-1}L(x, t, D)u - u_t = f(x, t), (x, t) \in \Omega_{\infty},$$
 (4)

such that

$$u(x, 0) = 0$$
 on  $\Omega$ , (5)

$$\frac{\partial^{j} u}{\partial \nu^{j}}|_{\Gamma_{\infty}} = 0, \quad j = 0, \dots, m - 1, \tag{6}$$

where  $\nu$  is the unit outward normal to  $\partial\Omega$ .

**Definition.** A function u(x,t) is called a generalized solution of (4)-(6) in  $H^{m,1}(\Omega_{\infty})$  if  $u(x,t) \in \mathring{H}^{m,1}(\Omega_{\infty})$ , u(x,0) = 0 and for all T > 0, u(x,t) satisfies the following integnal identity on  $\Omega_T$ :

$$(-1)^{m-1}\int\limits_{\Omega_T}\left[\sum_{|p|,|q|=1}^m(-1)^{|p|}a_{pq}D^qu\overline{D^p\eta}+\sum_{|p|=1}^ma_pD^pu\overline{\eta}+au\overline{\eta}\right]dxdt-$$

$$-\int_{\Omega_T} u_t \bar{\eta} dx dt = \int_{\Omega_T} f \bar{\eta} dx dt \quad \forall \eta \in \mathring{\hat{H}}^{m,0}(\Omega_T), \quad \eta(x,T) = 0. \tag{7}$$

Theorem 1. If  $a_{pq}$  are continuous on  $\overline{\Omega}_{\infty}$  whenever |p| = |q| = m and

$$\left| \frac{\partial a_{pq}}{\partial t}, \frac{\partial a_p}{\partial t} \right| \le \mu, \quad 1 \le |p|, \quad |q| \le m, \quad \mu = \text{const},$$

then problem (4)-(6) has at most one generalized solution in  $H^{m,1}(\Omega_{\infty})$ .

**Proof.** Suppose that problem (4)-(6) has two solutions  $u_1$ ,  $u_2$  and denote  $u(x,t)=u_1(x,t)-u_2(x,t)$ . For any T>0 and  $b\in (0,T)$ , we define  $\eta(x,t)=\int_b^t u(x,\tau)d\tau$  for  $0\leq t\leq b$  and  $\eta(x,t)=0$  for  $b\leq t< T$ . One can see that

$$\eta_t(x,t) = u(x,t), \quad 0 \le t \le b, \qquad \eta(x,t) \in \mathring{H}^{m,0}(\Omega_T)$$

and by identity (7), we have

$$(-1)^{m-1}\int\limits_{\Omega_b}\left[\sum_{|p|,|q|=1}^m(-1)^{|p|}a_{pq}D^q\eta_t\overline{D^p\eta}+\sum_{|p|=1}^ma_pD^p\eta_t\overline{\eta}\right]dxdt+\\$$

$$+(-1)^{m-1}\int_{\Omega_b}a\eta_t\overline{\eta}dxdt-\int_{\Omega_b}\eta_{tt}\overline{\eta}dxdt=0. \tag{8}$$

Let  $a_1 = a - (-1)^m \lambda_1 I$ , where  $\lambda_1$  is a positive number satisfying (3). Thus,

$$(-1)^{m-1} \int_{\Omega_b} \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} a_{pq} D^q \eta_t \overline{D^p \eta} + (-1)^m \lambda_1 \eta_t \overline{\eta} + \right]$$

$$+\sum_{|p|=1}^{m} a_{p} D^{p} \eta_{t} \overline{\eta} + a_{1} \eta_{t} \overline{\eta} \right] dx dt - \int_{\Omega_{b}} \eta_{tt} \overline{\eta} dx dt = 0.$$
 (9)

Thus, the integration of the real parts of (9) gives

$$2\int_{\Omega_b} |\eta_t|^2 dx dt + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \int_{\Omega} a_{pq}(x,0) D^q \eta(x,0) \overline{D^p \eta(x,0)} dx + (-1)^m \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} \right] dx + (-1)$$

$$\begin{split} &+2\mathrm{Re}\sum_{|p|=1}^{m}\int\limits_{\Omega}a_{p}(x,0)D^{p}\eta(x,0)\overline{\eta(x,0)}dx\bigg] +\lambda_{1}\int\limits_{\Omega}|\eta(x,0)|^{2}dx = \\ &=(-1)^{m-1}\int\limits_{\Omega_{b}}\sum_{|p|,|q|=1}^{m}(-1)^{|p|}\frac{\partial a_{pq}}{\partial t}D^{q}\eta\overline{D^{p}\eta}dxdt + \\ &+(-1)^{m-1}2\mathrm{Re}\int\limits_{\Omega_{b}}\bigg\{\sum_{|p|=1}^{m}\bigg[\frac{\partial a_{p}}{\partial t}D^{p}\eta\overline{\eta}+a_{p}D^{p}\eta\overline{\eta_{t}}\bigg] +a_{1}\eta_{t}\overline{\eta}\bigg\}dxdt. \end{split}$$

Since  $\frac{\partial a_{pq}}{\partial t}$ ,  $\frac{\partial a_p}{\partial t}$  are bounded, by Cauchy inequality and (3) we obtain

$$\|\eta_t\|_{L_2(\Omega_b)}^2 + \|\eta(x,0)\|_{H^m(\Omega)}^2 \le C(\varepsilon) \sum_{|p|=0}^m \int_{\Omega_b} |D^p \eta|^2 dx dt + \varepsilon \|\eta_t\|_{L_2(\Omega_b)}^2$$
 (10)

and

$$\|\eta(x,0)\|_{H^m(\Omega)}^2 \le C \sum_{|p|=0}^m \int_{\Omega_k} |D^p \eta|^2 dx dt.$$
 (11)

Setting

$$v_p(x,t) = \int_t^0 D^p u(x,\tau) d\tau, \quad 0 < t < b, \qquad J(t) = \sum_{|p|=0}^m \int_{\Omega} |v_p(x,t)|^2 dx,$$

we have

$$(1 - Cb)J(b) \le C \int_{0}^{b} J(t)dt, \quad b \in \left[0, \frac{1}{2}C\right]. \tag{12}$$

By Gronwall-Bellman inequality, one has  $J(t)\equiv 0$ . Thus, (10), (11) lead to  $\eta_t\equiv 0$ , i.e.,  $u_1\equiv u_2 \ \forall t\in [0,1/2C]$ . Using the same argument as before for functions  $u_1$ ,  $u_2$  on [1/2C,T], we can show that  $u_1\equiv u_2 \ \forall t\in [1/2C,1/C]$ . Continuing in this fashion, after finite number of steps, we can prove that  $u_1\equiv u_2 \ \forall t\in [0,T]$ . Since T>0 is arbitrary,  $u_1\equiv u_2 \ \forall t\in [0,\infty)$ . This completes the proof.

Now we prove the following theorem as a result on existence.

Theorem 2. Suppose that  $a_{pq}$  are continuous on  $\overline{\Omega}_{\infty}$  whenever |p| = |q| = m, there exists a  $\lambda > 0$ , a function  $B(t) \in L_1(0, +\infty)$  such that  $fe^{\lambda t} \in L_2(\Omega_{\infty})$  and

$$\max\left\{\left|\frac{\partial a_{pq}}{\partial t}\right|, |a_p|, |a-(-1)^m(\lambda+\lambda_0)I|\right\} \leq B(t) \quad \forall (t,x) \in \Omega_{\infty},$$

where  $\lambda_0$  is determined from (2). Then problem (4)-(6) has a generalized solution  $u(x,t) \in H^{m,1}(\Omega_{\infty})$  satisfying

$$||u||_{H^{m,1}(\Omega_{\infty})}^2 \le C||fe^{\lambda t}||_{L_2(\Omega_{\infty})}^2.$$

**Proof.** Suppose that  $\{\varphi_k(x)\}_{k=1}^{\infty} \subset \overset{\circ}{H}^m(\Omega)$  is orthonormal in  $L_2(\Omega)$  and its linear closure is  $\overset{\circ}{H}^m(\Omega)$ . For each natural number N, let us consider the function

$$u^{N}(x, t) = \sum_{k=1}^{N} C_{k}^{N}(t)\varphi_{k}(x),$$

where  $C_k^N(t)$  satisfies  $C_k^N(0) = 0$  and

$$\int_{\Omega} \left[ u_t^N \overline{\varphi_\ell} + \sum_{|p|,|q|=1}^m (-1)^{m+|p|} a_{pq} D^q u^N \overline{D^p \varphi_\ell} + \lambda_0 u^N \overline{\varphi_\ell} \right] dx + \\
+ (-1)^m \int_{\Omega} \left( \sum_{|p|=1}^m a_p D^p u^N + a_0 u^N \right) \overline{\varphi_\ell} dx = \\
= - \int_{\Omega} f \overline{\varphi_\ell} dx, \qquad \ell = 1, \dots, N, \tag{13}$$

with  $a_0 = a - (-1)^m \lambda_0 I$ . Put  $v^N(x,t) = u^N(x,t)e^{\lambda t}$ . Multiplying (13) by  $e^{\lambda t} \frac{d}{dt} (\overline{C_\ell^N(t)e^{\lambda t}})$ , then taking the sum in  $\ell$  from 1 to N, one obtains

$$\int_{\Omega} |v_t^N|^2 dx + \int_{\Omega} \sum_{|p|,|q|=1}^m \left[ (-1)^{m+|p|} a_{pq} D^q v^N \overline{D^p v_t^N} + \lambda_0 v^N \overline{v_t^N} \right] dx +$$

$$+ (-1)^m \int_{\Omega} \left[ \sum_{|p|=1}^m a_p D^p v^N \overline{v_t^N} + (a_0 - (-1)^m \lambda I) v^N \overline{v_t^N} \right] dx = - \int_{\Omega} f e^{\lambda t} \overline{v_t^N} dx. \tag{14}$$

Since  $a_{pq} = (-1)^{|p|+|q|} a_{qp}^*$ , we have

$$2\operatorname{Re}\left[\sum_{|p|,|q|=1}^{m}(-1)^{|p|}a_{pq}D^{q}v^{N}\overline{D^{p}v_{t}^{N}}+(-1)^{m}\lambda_{0}v^{N}\overline{v_{t}^{N}}\right]=$$

$$=\frac{\partial}{\partial t}\left(\sum_{|p|,|q|=1}^{m}(-1)^{|p|}a_{pq}D^{q}v^{N}\overline{D^{p}v^{N}}+(-1)^{m}\lambda_{0}|v^{N}|^{2}\right)-$$

$$-\sum_{|p|,|q|=1}^{m}(-1)^{|p|}\frac{\partial a_{pq}}{\partial t}D^{q}v^{N}\overline{D^{p}v^{N}}.$$
(15)

By integrating the real parts of (14) in t from 0 to t and using (15), we get

$$2\int\limits_{\Omega_t}|v_t^N|^2dxdt+$$

$$+ \int_{\Omega} \sum_{|p|,|q|=1}^{m} \left[ (-1)^{m+|p|} a_{pq}(x,t) D^{q} v^{N}(x,t) \overline{D^{p} v^{N}(x,t)} + \lambda_{0} |v^{N}(x,t)|^{2} \right] dx =$$

$$= \int_{\Omega_{t}} \sum_{|p|,|q|=1}^{m} (-1)^{m+|p|} \frac{\partial a_{pq}}{\partial t} D^{q} v^{N} \overline{D^{p} v^{N}} dx dt -$$

$$-(-1)^m 2 \mathrm{Re} \int\limits_{\Omega_t} \Bigg[ \sum_{|p|=1}^m a_p D^p v^N \overline{v_t^N} + \overline{a} v^N \overline{v_t^N} \Bigg] dx dt - 2 \mathrm{Re} \int_{\Omega_t} f e^{\lambda t} \overline{v_t^N} dx dt,$$

where  $\overline{a} = (a_0 - (-1)^m \lambda I)$ .

By (2) and the Cauchy inequality, we obtain

$$2\int\limits_{\Omega_{t}}|v_{t}^{N}|^{2}dxdt+\mu_{0}||v^{N}(x,t)||_{H^{m}(\Omega)}^{2}\leq$$

$$\leq C_1(\varepsilon)\int\limits_0^t B(t)\|v^N(x,t)\|_{H^m(\Omega)}^2 dt +$$

$$+ \varepsilon \int\limits_{\Omega_t} |v_t^N|^2 dx dt + C_2(\varepsilon) \int\limits_0^t \|fe^{\lambda t}\|_{L_2(\Omega)}^2 dt.$$

Therefore,

$$\int\limits_{\Omega_t}|v^N_t|^2dxdt+\|v^N(x,t)\|^2_{H^m(\Omega)}\leq$$

$$\leq C_3 \int_{0}^{t} B(t) \|v^{N}(x,t)\|_{H^{m}(\Omega)}^{2} dt + C_4 \int_{0}^{\infty} \|fe^{\lambda t}\|_{L_2(\Omega)}^{2} dt.$$
 (16)

So we get

$$\|v^N(x,t)\|_{H^m(\Omega)}^2 \le C_3 \int_0^t B(t) \|v^N(x,t)\|_{H^m(\Omega)}^2 dt + C_4 \|fe^{\lambda t}\|_{L_2(\Omega_\infty)}^2.$$

By Gronwall-Bellman inequality, we have

$$||v^{N}(x,t)||_{H^{m_{i}}(\Omega)}^{2} \le C_{4}||fe^{\lambda t}||_{L_{2}(\Omega_{\infty})}^{2}e^{C_{3}\int_{0}^{t}B(t)dt}$$
. (17)

Thus,

$$||v^N(x,t)||^2_{H^m(\Omega)} \le C_5 ||fe^{\lambda t}||^2_{L_2(\Omega_{sol})}$$

i.e.,

$$e^{2\lambda t} \|u^N(x,t)\|_{H^m(\Omega)}^2 \le C_5 \|fe^{\lambda t}\|_{L_2(\Omega_\infty)}^2$$

Hence,

$$||u^{N}(x,t)||_{H^{m,0}(\Omega_{\infty})}^{2} \le C_{6}||fe^{\lambda t}||_{L_{2}(\Omega_{\infty})}^{2}$$
 (18)

Multiplying (13) by  $\frac{d(C_{\ell}^{N}(t))}{dt}$ , taking the sum in  $\ell$  from 1 to N, and integrating in t from 0 to t, we obtain

$$\int\limits_{\Omega_t} \left[ u_t^N \overline{u_t^N} + \sum_{|p|,|q|=1}^m (-1)^{m+|p|} a_{pq} D^q u^N \overline{D^p u_t^N} + \lambda_0 u^N \overline{u_t^N} \right] dx dt +$$

$$+(-1)^{m}\int\limits_{\Omega_{t}}\left(\sum_{|p|=1}^{m}a_{p}D^{p}u^{N}\overline{u_{t}^{N}}+a_{0}u^{N}\overline{u_{t}^{N}}\right)dxdt=-\int\limits_{\Omega_{t}}f\overline{u_{t}^{N}}dxdt. \tag{19}$$

Hence,

$$2||u_t^N||_{L_2(\Omega_t)}^2 +$$

$$+\int\limits_{\Omega}\Bigg[\sum_{|p|,|q|=1}^{m}(-1)^{m+|p|}a_{pq}(x,t)D^{q}u^{N}(x,t)\overline{D^{p}u^{N}(x,t)}+\lambda_{0}\|u^{N}(x,t)\|^{2}\Bigg]dx=$$

$$=\int\limits_{\Omega_L}\sum\limits_{|p|,|q|=1}^m(-1)^{m+|p|}\frac{\partial a_{pq}}{\partial t}D^qu^N\overline{D^pu^N}\,dx\,dt-$$

$$-(-1)^m 2 \operatorname{Re} \int\limits_{\Omega_t} \left[ \sum_{|p|=1}^m a_p D^p u^N \overline{u_t^N} + a_0 u^N \overline{u_t^N} \right] dx dt - 2 \operatorname{Re} \int\limits_{\Omega_t} f \overline{u_t^N} dx dt.$$

By Cauchy inequality and (2) we obtain

$$||u_t^N||_{L_2(\Omega_t)}^2 + ||u^N(x,t)||_{H^m(\Omega)}^2 \le D_1 ||u^N||_{H^{m,0}(\Omega_t)}^2 + D_2 ||f||_{L_2(\Omega_t)}^2$$

Thus,

$$||u_{t}^{N}||_{L_{2}(\Omega_{t})}^{2} \leq D_{1}||u^{N}||_{H^{m,0}(\Omega_{t})}^{2} + D_{2}||f||_{L_{2}(\Omega_{t})}^{2} \Rightarrow ||u_{t}^{N}||_{L_{2}(\Omega_{\infty})}^{2} \leq$$

$$\leq D_{1}||u^{N}||_{H^{m,0}(\Omega_{\infty})}^{2} + D_{2}||f||_{L_{2}(\Omega_{\infty})}^{2} \leq$$

$$\leq D_{1}||u^{N}||_{H^{m,0}(\Omega_{\infty})}^{2} + D_{2}||fe^{\lambda t}||_{L_{2}(\Omega_{\infty})}^{2}.$$
(20)

According to (18), we obtain

$$||u_t^N||_{L_2(\Omega_\infty)}^2 \le (C_6D_1 + D_2)||fe^{\lambda t}||_{L_2(\Omega_\infty)}^2.$$
 (21)

Thus, (18), (21) imply

$$||u^N||_{H^{m,1}(\Omega_\infty)}^2 \le C||fe^{\lambda t}||_{L_2(\Omega_\infty)}^2,$$
 (22)

where C does not depend on N.

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Since the sequence of functions  $\{u^N\}$  is uniformly bounded in  $H^{m,1}(\Omega_\infty)$ , we can take a subsequence which is weakly convergent in  $H^{m,1}(\Omega_\infty)$  to some function u(x,t). We will prove that u(x,t) is the solution of (4)-(6). Since  $u^N(x,0)=0$   $\forall x\in\Omega$  and  $u^N(x,t)\in \mathring{H}^{m,1}(\Omega_T)$ , it follows that  $u(x,0)=0, u(x,t)\in \mathring{H}^{m,1}(\Omega_T)$ . For any T>0, multiplying (13) by  $d_\ell(t)\in H^1(0,T),\ d_\ell(T)=0$ , taking the sum in  $\ell$  from 1 to N and integrating in t from 0 to t, we obtain

$$\int_{\Omega_{T}} u_{t}^{N} \overline{\eta} dx dt +$$

$$+(-1)^{m} \int_{\Omega_{T}} \left[ \sum_{|p|,|q|=1}^{m} (-1)^{|p|} a_{pq} D^{q} u^{N} \overline{D^{p} \eta} + (-1)^{m} u^{N} \overline{\eta} \right] dx dt +$$

$$+(-1)^{m} \int_{\Omega_{T}} \left[ \sum_{|p|=1}^{m} a_{p} D^{p} u^{N} \overline{\eta} + a_{0} u^{N} \overline{\eta} \right] dx dt =$$

$$= - \int_{\Omega_{T}} f \overline{\eta} dx dt. \qquad (23)$$

The above equality is true for any function  $\eta \in M_N$ , where  $M_N$  is the set consisting of functions which have the form  $\sum_{i=1}^N d_i(t)\varphi_i(x)$ ,  $d_i(t) \in H^1(0,T)$ ,  $d_i(T)=0$ . Since  $\{u^N\}$  is weakly convergent, passing to the limit as  $N \to \infty$  in equality (23). We get

$$\int_{\Omega_T} u_t \overline{\eta} dx dt + (-1)^m \int_{\Omega_T} \left[ \sum_{|p|,|q|=1}^m (-1)^{|p|} a_{pq} D^q u \overline{D^p \eta} + (-1)^m u \overline{\eta} \right] dx dt + \\
+ (-1)^m \int_{\Omega_T} \left[ \sum_{|p|=1}^m a_p D^p u \overline{\eta} + a_0 u \overline{\eta} \right] dx dt = - \int_{\Omega_T} f \overline{\eta} dx dt. \tag{24}$$

Since  $M = \bigcup_{N=1}^{\infty} M_N$  is dense in  $\hat{H}^{m,0}(\Omega_T)$ , it follows that (24) holds for any function  $\eta \in \mathring{H}^{m,0}(\Omega_T)$ ,  $\eta(x,T) = 0$ , i.e., u(x,t) is a generalized solution of (4)–(6). Moreover, the weak convergence of  $\{u^N\}$  and (22) imply that

$$||u(x,t)||_{H^{m,1}(\Omega_{\infty})} = \underline{\lim}_{N \to \infty} ||u^N(x,t)||_{H^{m,1}(\Omega_{\infty})} \le C||fe^{\lambda t}||_{L_2(\Omega_{\infty})}.$$

The theorem is proved.

3. Asymptotic behaviour of solutions. We now consider the asymptotic behaviour as  $t \to \infty$  of generalized solutions of the first boundary-value problem

$$(-1)^{m-1}L(x,t,D)u - u_t = f(x,t), (25)$$

where  $(x,t) \in \Omega_{\infty}$ , such that

$$u(x,t)|_{t=0} = \varphi(x) \in \mathring{H}^{2m}(\Omega),$$
 (26)

$$\frac{\partial^{j} u}{\partial \nu^{j}}\Big|_{\Gamma_{\infty}} = 0, \quad j = 0, \dots, m-1.$$
 (27)

Theorem 3. Suppose that

1)  $a_{pq}$  are continuous on  $\overline{\Omega}_{\infty}$  whenever |p| = |q| = m;

2) 
$$\left| \frac{\partial a_{pq}}{\partial t}, \frac{\partial a_{p}}{\partial t} \right| \le \mu, \ 1 \le |p|, \ |q| \le m, \ \mu = \text{const};$$

there exist λ<sub>1</sub> > λ<sub>2</sub> > 0 and a function B(t) ∈ L<sub>1</sub>(0,+∞) such that fe<sup>λ<sub>1</sub>t</sup> ∈ E<sub>2</sub>(0,+∞), and

$$\max\left\{\left|\frac{\partial a_{pq}}{\partial t}\right|,\left|\frac{\partial a_p}{\partial t}\right|,|a-(-1)^m(\lambda_0+\lambda_2)I|\right\}\leq B(t)\quad\forall (t,x)\in\Omega_\infty.$$

Then problem (25)-(27) has a generalized solution u(x,t) in  $H^{m,1}(\Omega_{\infty})$ . Moreover, we have

$$u(t,x) = \varphi(x)e^{-\lambda_1 t} + o(e^{-\frac{\lambda_2}{2}t})$$
 as  $t \to \infty$ .

**Proof.** We set  $v(x,t) = u(x,t) - \varphi(x)e^{-\lambda_1 t}$ . Then problem (25)-(27) can be written as

$$(-1)^{m-1}L(x,t,D)v - v_t = f(x,t) + e^{-\lambda_1 t} [(-1)^m L(x,t,D)\varphi - \lambda_1 \varphi(x)], \quad (28)$$

where  $(x,t) \in \Omega_{\infty}$ , such that

$$v(x,0) = 0$$
 on  $\Omega$ , (29)

$$\frac{\partial^{j} v}{\partial \nu^{j}}\Big|_{\Gamma_{\infty}} = 0, \quad j = 0, \dots, m-1.$$
 (30)

Consider  $g(x,t) = f(x,t) + e^{-\lambda_1 t}[(-1)^m L(x,t,D)\varphi - \lambda_1 \varphi(x)]$ . It is easy to see that  $ge^{\lambda_2 t} \in L_2(\Omega_\infty)$ . Thus, by virtue of Theorem 2, problem (28)-(30) has a generalized solution  $v(x,t) \in H^{m,1}(\Omega_\infty)$ . Furthermore, by the same proof for Theorem 2, as  $N \to \infty$ , inequality (16) implies that

$$||e^{\lambda_2 t}v(x,t)||^2_{H^m(\Omega)} \le C_1,$$
 (31)

$$\|(e^{\lambda_2 t}v(x,t))_t\|_{L_2(\Omega_{\infty})}^2 \le C_2.$$
 (32)

It follows from (31), (32) that

$$e^{\lambda_2 t} \Big[ \|v(x,t)\|_{L_2(\Omega)}^2 + \|Dv(x,t)\|_{L_2(\Omega)}^2 \Big] \leq C_1 e^{-\lambda_2 t},$$

$$\int\limits_{0}^{\infty}e^{\lambda_{2}t}\|v_{t}(x,t)\|_{L_{2}(\Omega)}^{2}dt\leq$$

$$\leq \int\limits_0^\infty e^{-\lambda_2 t} \|(e^{\lambda_2 t}v(x,t))_t\|_{L_2(\Omega)}^2 dt +$$

$$+ \int_{0}^{\infty} \lambda_{2}^{2} e^{\lambda_{2} t} \|v(x,t)\|_{L_{2}(\Omega)}^{2} dt \le C_{3}.$$

Thus,

$$\int_{0}^{\infty} e^{\lambda_{2}t} \Big[ \|v_{t}(x,t)\|_{L_{2}(\Omega)}^{2} + \|v(x,t)\|_{L_{2}(\Omega)}^{2} + \|Dv(x,t)\|_{L_{2}(\Omega)}^{2} \Big] dt < \infty.$$

By virtue of the Nash inequality [10], this gives

$$|v(x,t)| \le Ce^{-\frac{\lambda_2}{2}t}$$
 as  $t \to \infty$ .

This implies that the generalized solution u(x,t) of problem (25)-(27) satisfies the inequality

$$|u(x,t) - e^{-\lambda_1 t} \varphi(x)| \le C e^{-\frac{\lambda_2}{2} t}$$
.

The proof is completed.

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