

Nature of extrinsic and intrinsic self-trapping of charge carriers in underdoped cuprate high- T_c superconductors

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Nature of extrinsic and intrinsic self-trapping (ST) of charge carriers in cuprates have been studied theoretically. The binding energies and radii of the extrinsic and intrinsic large polarons and bipolarons in cuprates are calculated variationally using the continuum model and adiabatic approximation. We have shown that the extrinsic and intrinsic three-dimensional (3D) large bipolarons exist in underdoped cuprates at $\eta = \varepsilon_\infty/\varepsilon_0 < 0.127$ and $\eta < 0.138$, respectively [where ε_∞ (ε_0) is the optic (static) dielectric constant].

Key words: polaron, bipolaron, self-trapping, high- T_c superconductors

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1. Introduction

It is well known that electron-phonon coupling (EPC) is one of the most common and fundamental interactions in solids. In particular, a strong EPC in high-temperature superconducting cuprates (HTSC) was reported by a lot of experiments, which implies that EPC plays an important role in pairing. Accordingly, charge carriers in polar materials interacting with the phonon field can cause the formation of self-trapped polarons and bipolarons. The interest to the polarons and bipolarons is caused by their important role in explaining many characteristics of HTSC compounds (see [1, 2] and references therein). The ground state of an undoped HTSC is an antiferromagnet (AF) whose doping by holes leads to the phase showing high temperatures of the superconducting (SC) transition. The extended $t-t'-t''-J$ -Hubbard model, where a hole in a two-dimensional (2D) AF is simultaneously magnetic and a lattice polaron that moves through the lattice emitting and absorbing magnons and phonons, is used in many papers for calculation of spectral properties of underdoped HTSC compounds [3]. According to the band calculations and angle resolved-photoemission spectroscopy (ARPES) data [4, 5], the electronic structure of the parent cuprate compounds is well described by a three-band Hubbard model and the oxygen valence band lies within the Mott-Hubbard gap. Furthermore, according to the combined study of the ARPES and quantum Monte Carlo simulations, there was proposed a novel polaronic metallic state in underdoped cuprates [3]. The variety of opinions range from a complete negation of the role of EPC in the physics of HTSC [6, 7] to the statement that the binding energy of the polaron is an order of magnitude larger than any characteristic energy of the magnetic subsystem, and, therefore, exactly the magnetic system is irrelevant [8, 9]. Both of the above radical statements have been criticized many times. On the other hand, the success of the extended $t-J$ -Hubbard model in describing the ARPES [3, 10] and optical spectra [11] does not make it possible to completely disregard the role of the magnetic subsystem. Correspondingly, there is an opinion that interactions with both magnetic and lattice subsystems are important [12]. Results indicating an important role of the EPC in HTSC compounds are presented, with emphasis on its implications for ARPES and optical conductivity [12].

Hole doping of the cuprates produces first quasi-free holes having the mass m_h in the oxygen valence band. The hole carriers are assumed to be within both a three-dimensional (3D) and a 2D deformable

medium, the last one being CuO_2 layers [13]. In reality, however, no systems can be purely 2D, and therefore, the layered cuprate compounds may be approximated as a 3D deformable medium. There is also a convincing experimental evidence that the consideration of cuprates as 3D systems may appear to be more appropriate (see [14–17]). The experimental results presented in [16] indeed confirm that the hole-doped system $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ (LSCO) becomes less 2D in the strongly localized state. In polar materials, the hole carriers interacting both with lattice vibrations (i.e., acoustic and optical phonons) and with lattice defects (e.g., dopants or impurities), can easily be self-trapped near the defects and in a defect-free deformable lattice. Actually, the carrier localization in the cuprates is interrelated and the quantitative theory of this phenomenon is still lacking. Particularly, possible roles of large- and small-radius dopants, dopant-driven and carrier-driven inhomogeneities, carrier–defect–lattice and carrier–lattice interactions and other factors are very important for the localization of carriers in hole-doped cuprates. In the present paper we study the extrinsic (defect-assisted) and intrinsic (phonon-assisted) single particle and pair self-trapping (ST) of carriers using the continuum model of ionic crystal and adiabatic approximation. The possibility of the formation of localized extrinsic and intrinsic (bi)polaronic states as well as hydrogen-like impurity states in charge-transfer (CT) gap of the cuprates is examined and compared with the experimental data.

2. Calculation of the ground-state energy of the system of a defect-bound hole carrier in the polar crystal

Electron-phonon interaction (EPI) affects the electronic properties of semiconductors and polar crystals in various ways depending on the strength of the electron-phonon interaction. Among them, the polaron formation and the ST are its typical and important effects on the carriers. The relevant charge carriers in hole-doped cuprates are large polarons [13, 18] and the strong EPIs are responsible for enhancement of the polaron mass $m_p = (2.0\text{--}3.0)m_h$ [19] (where $m_h = m_e$ is the free electron mass). According to Toyozawa [20], the mechanisms for ST of carriers are classified as intrinsic and extrinsic ones. The intrinsic mechanism means that carriers are self-trapped at deformed lattice sites through EPI. The extrinsic mechanism consists of short-range and/or long-range potentials by impurities or defects enhancing ST of carriers due to EPI. So far, there are no detailed quantitative studies of the intrinsic and extrinsic ST of the carriers in lightly doped cuprates. In order to better understand the situation, the possibility of the formation of extrinsic self-trapped states and of the intrinsic ones in the CT gap of the cuprates need to be thoroughly examined and compared with the existing experimental results confirming the existence of such localized in-gap states.

We use a continuum model proposed in [20] and adiabatic approximation to calculate the ground-state energy of an interacting system of a defect (dopant)-bound hole carrier in the polar crystal. In the continuous approximation, the functionals of the total energies of the single-carrier and two-carrier systems can be written as follows:

$$E_1[\psi(r)] = \frac{\hbar^2}{2m^*} \int |\nabla\psi(r)|^2 d^3r - \frac{e^2}{2\bar{\epsilon}} \int \frac{\psi^2(r)\psi^2(r')}{|r-r'|} d^3r d^3r' - \frac{E_d^2}{2K} \int \psi^4(r) d^3r - \frac{Ze^2}{\epsilon_0} \int \frac{\psi^2(r)}{r} d^3r + \left(V_0 - \frac{E_d E_{dD}}{K} \right) \int \psi^2(r) \delta(r) d^3r \quad (2.1)$$

and

$$E_2[\Psi(r_1, r_2)] = \frac{\hbar^2}{2m^*} \int [|\nabla_1\Psi(r_1, r_2)|^2 + |\nabla_2\Psi(r_1, r_2)|^2] d^3r_1 d^3r_2 + \frac{e^2}{\epsilon_\infty} \int \frac{\Psi^2(r_1, r_2)}{|r_1 - r_2|} d^3r_1 d^3r_2 - \frac{2e^2}{\bar{\epsilon}} \int \frac{\Psi^2(r_1, r_2)\Psi^2(r_3, r_4)}{|r_1 - r_3|} d^3r_1 d^3r_2 d^3r_3 d^3r_4 - \frac{2E_d^2}{K} \int \Psi^2(r_1, r_2)\Psi^2(r_2, r_3) d^3r_1 d^3r_2 d^3r_3 - \frac{2Ze^2}{\epsilon_0} \int \frac{\Psi^2(r_1, r_2)}{r_1} d^3r_1 d^3r_2 + 2 \left(V_0 - \frac{E_d E_{dD}}{K} \right) \int \Psi^2(r_1, r_2) \delta(r_1) d^3r_1 d^3r_2, \quad (2.2)$$

where m^* and e are carrier effective mass and charge, respectively, r_1 and r_2 are the position vectors of the carriers, E_d and E_{dD} are the deformation potentials of the carrier and the defect, respectively, $\tilde{\epsilon}$ is given by $\tilde{\epsilon}^{-1} = \epsilon_\infty^{-1} - \epsilon_0^{-1} = (1 - \eta)/\epsilon_\infty$, K is an elastic constant, V_0 is the short-range defect potential, Z is the charge state of the defect. A large ionicity of the cuprates $\eta = \epsilon_\infty/\epsilon_0 \ll 1$, enhances the polar EPI and the tendency to polaron formation.

In order to minimize the functionals (2.1) and (2.2) with respect to $\psi(r)$ and $\Psi(r_1, r_2)$ we can take the trial wave functions as follows:

$$\psi(r) = N_1 \exp[-(\sigma r)] \quad (2.3)$$

and

$$\Psi(r_1, r_2) = N_2 [1 + \gamma(\sigma r_{12})] \exp[-\sigma(r_1 + r_2)], \quad (2.4)$$

where $N_1 = \sigma^{3/2}/\sqrt{\pi}$ and $N_2 = \sigma^3/\pi \sqrt{K_1(\gamma)}$ are the normalization factors, $\sigma = \beta/a_0$, $K_1(\gamma) = 1 + \frac{35}{8}\gamma + 6\gamma^2$ is the correlation coefficient, β and γ are the variational parameters that characterize the localization degree of carriers and the correlation in their motions, respectively. $r_{12} = |r_1 - r_2|$ is the distance between the carriers, a_0 is the lattice constant. Substituting equations (2.3) and (2.4) into equations (2.1) and (2.2), and performing the integrations in equations (2.1) and (2.2), we obtain the following functionals

$$E_1[\beta] = A \left[\beta^2 - g_s(1 + b_s)\beta^3 - g_l \left(1 - \eta + \frac{16}{5} Z \eta \right) \beta \right] \quad (2.5)$$

and

$$\begin{aligned} E_2[\beta, \gamma] = & 2A \frac{K_2(\gamma)}{K_1(\gamma)} \left\{ \beta^2 - \frac{8}{5} g_l \left[2(1 - \eta) \frac{K_4(\gamma)}{K_1(\gamma)K_2(\gamma)} + Z \eta \frac{K_5(\gamma)}{K_2(\gamma)} - \frac{K_3(\gamma)}{K_2(\gamma)} \right] \beta \right. \\ & \left. - 16g_s \left[\frac{K_6(\gamma)}{K_1(\gamma)K_2(\gamma)} + \frac{b_s K_7(\gamma)}{16 K_2(\gamma)} \right] \beta^3 \right\}, \end{aligned} \quad (2.6)$$

where $A = \hbar^2/2m^* a_0^2$, $g_s = E_d^2/16\pi K a_0^3 A$ and $g_l = 5e^2/16\epsilon_\infty a_0 A$ are dimensionless short-range and long-range carrier-phonon coupling parameters, $b_s = 16(E_{dD}/E_d - K V_0/E_d^2)$ is short-range carrier-defect-phonon coupling parameter, and the correlation coefficients are analytical functions of the variational parameter γ :

$$\begin{aligned} K_2(\gamma) &= 1 + \frac{25}{8}\gamma + 4\gamma^2, \\ K_4(\gamma) &= \frac{5}{8} + \frac{1087}{216}\gamma + \frac{38237}{2304}\gamma^2 + \frac{67639}{2592}\gamma^3 + \frac{4293}{256}\gamma^4, \\ K_3(\gamma) &= \frac{5}{8}\gamma + 2\gamma + \frac{35}{16}\gamma^2, \quad K_5(\gamma) = 2 + \frac{15}{2}\gamma + 9\gamma^2, \\ K_6(\gamma) &= \frac{1}{8} + \frac{185}{216}\gamma + \frac{4199}{1728}\gamma^2 + \frac{8591}{2592}\gamma^3 + \frac{477}{256}\gamma^4, \\ K_7(\gamma) &= 1 + 3\gamma + 3\gamma^2. \end{aligned}$$

Minimization of the functionals (2.5) and (2.6) over the variational parameters β and γ would give the ground state energies $E_1^{\min}(\beta)$ and $E_2^{\min}(\beta, \gamma)$ of single and pair self-trapped carriers. As a function of β and γ , these functionals have also got the maxima $E_1^{\max}(\beta)$ and $E_2^{\max}(\beta, \gamma)$. The other parameters A , g_s , g_l , b_s and Z entering equations (2.5) and (2.6) play different roles in the formation of self-trapped states and determine the nature of the localized states of carriers in doped polar materials. The values of the parameters A , g_s and g_l can be obtained using the experimental values of the parameters a_0 , m^* , K , ϵ_∞ , ϵ_0 and the Fermi energy E_F of the undoped cuprates. The parameters b_s and Z characterize the formation of extrinsic (i.e., hydrogenic and non-hydrogenic) self-trapped states of carriers, whereas the parameters g_s and g_l characterizing the strengths of short- and long-range carrier-phonon interactions are responsible for the formation of intrinsic self-trapped states.

Using the equations (2.5) and (2.6), we calculate the energies of different localized in-gap states in the CT gap of the cuprates. In order to determine the nature of these in-gap states and the quasi-free to localized state transition, we distinguish different physical situations in these systems. One can make

interesting analyses of real systems based on the sign and magnitude of b_s . In hole-doped cuprates, the situations might be quite different for different types of dopants. In particular, the signs of the deformation potential constants E_d and E_{dD} for holes and small-radius defects are always positive, while E_{dD} for large-radius defects is negative [20]. However, at present no information is available for the magnitudes of the parameters E_{dD} and V_0 . Therefore, the parameter b_s in equations (2.5) and (2.6) can be considered as the free parameter. We consider first the possibility of the formation of localized in-gap states at single and pair ST of carriers near the small-radius dopants (with $E_{dD} > 0$ or $b_s > 0$) in La-based cuprates. In this case, both short and long range parts of the defect potential in equations (2.1) and (2.2) are attractive, so that the substitution of small-radius cations (e.g., Ca^{2+} and Nd^{3+} ions) for La^{3+} ions in La_2CuO_4 and for Sr^{2+} ions in LSCO leads to a combined defect- and phonon-assisted ST of hole carriers with the formation of localized single-carrier and two-carrier impurity states, which are extrinsic polaronic and bipolaronic (the so-called U-pairing) states. At $Z \neq 0$ and $b_s > 0$, the minima of $E_1(\beta)$ and $E_2(\beta, \gamma)$ correspond to the ground-state energies of the extrinsic large polaron and bipolaron, respectively, measured with respect to the top of the oxygen valence band. The binding energies of such extrinsic large polaron and bipolaron (or negative U center) are defined as $E_{pI} = |E_1^{\min}(\beta)|$ and $E_{bU} = |E_2^{\min}(\beta, \gamma) - 2E_1^{\min}(\beta)|$, respectively. In 3D systems, there is generally a potential barrier between the large- and small-radius self-trapped states. The two states of the extrinsic large polaron are separated by a potential barrier, with activation energy $E_1^A = E_1^{\max}(\beta) - E_1^{\min}(\beta)$ needed for the transition from the large-radius localized state to the small-radius one. The potential barrier $E_2^A = E_2^{\max}(\beta, \gamma) - E_2^{\min}(\beta)$ exists between the large and small-radius extrinsic bipolaronic states.

We now calculate the basic parameters of the extrinsic large (bi)polarens in La-based cuprates. At low temperature, the La-based cuprates are orthorhombic with the lattice parameter $a_0 \simeq 5.4 \text{ \AA}$. According to the spectroscopy data, the Fermi energy of the undoped cuprates is about $E_F \simeq 7 \text{ eV}$ [21]. To determine the value of the short-range carrier-phonon coupling constant g_s , we can estimate the deformation potential E_d as $E_d = (2/3)E_F$ [22]. For the cuprates, typical values of other parameters are $m^* = m_e$ [13], $\epsilon_\infty = 3-5$ [23, 24], $K = 1.4 \cdot 10^{12} \text{ dyn/cm}^2$ [25], and $Z = 1$. The calculated values of E_{pI} , E_{bU} , E_1^A and E_2^A for $b_s = 1$ and different values of ϵ_∞ and η are presented in table 1. From table 1 we can see that the potential barriers separating the large- and small-radius extrinsic (bi)polarens are rather high. These high potential barriers prevent the formation of small extrinsic (bi)polarens in 3D cuprates. The defect- and phonon-assisted ST of large polaron and large bipolaron in La-based cuprates are shown in figures 1 and 2, respectively.

Table 1. The calculated parameters of the extrinsic large polarons and bipolarons (with correlation between the pairing carriers) in 3D cuprates at $Z = 1$, $b_s = 1$ and different values of ϵ_∞ and η .

η	$\epsilon_\infty = 3.5$				$\epsilon_\infty = 4$				$\epsilon_\infty = 4.5$			
	E_{pI}, eV	E_{bU}, eV	E_1^A, eV	E_2^A, eV	E_{pI}, eV	E_{bU}, eV	E_1^A, eV	E_2^A, eV	E_{pI}, eV	E_{bU}, eV	E_1^A, eV	E_2^A, eV
0.00	0.1135	0.0610	5.3373	5.8611	0.0863	0.0456	5.6962	6.4752	0.0679	0.0354	5.9807	6.7958
0.02	0.1240	0.0525	5.2128	5.7649	0.0943	0.0391	5.5849	6.3881	0.0741	0.0303	5.8801	6.7090
0.04	0.1349	0.0434	5.0893	5.6693	0.1026	0.0321	5.4742	6.3013	0.0806	0.0247	5.7800	6.6224
0.06	0.1464	0.0336	4.9668	5.5743	0.1113	0.0246	5.3643	6.2150	0.0874	0.0188	5.6804	6.5359
0.08	0.1584	0.0231	4.8453	5.4799	0.1203	0.0166	5.2552	6.1292	0.0945	0.0125	5.5815	6.4495
0.10	0.1709	0.0120	4.7248	5.3861	0.1298	0.0081	5.1468	6.0438	0.1019	0.0058	5.4831	6.3633
0.12	0.1839	0.0001	4.6053	5.2929	0.1396	—	5.0391	5.9588	0.1096	—	5.3854	6.2773

Another interesting question is in what way large-radius dopants in cuprates affect the carrier-phonon system, especially near such defects. This opposite situation is realized in LSCO or $\text{La}_{2-x}\text{Ba}_x\text{CuO}_4$ (LBCO), where the radius of Sr^{2+} ions is larger than that of La^{3+} ions [26], so that for Sr^{2+} ion $Z = 1$, $E_{dD} < 0$ or $b_s < 0$. In this case, the short-range part of the impurity potential in equation (2.1) is repulsive. Therefore, one can treat it like a hard core. The hole-lattice interactions near the large-radius dopants in LSCO and LBCO are suppressed by this repulsive defect potential and hole carriers are localized at a distance from the dopants (i.e., hole-carriers are loosely bound to dopants by long-range Coulomb attraction). From these considerations, it follows that the hole-lattice interaction near the large-radius dopants is weak and the localized impurity state may be of a hydrogen-like character described by a rigid lattice model [27]. Therefore, we can consider the hydrogen-like impurity centers having the

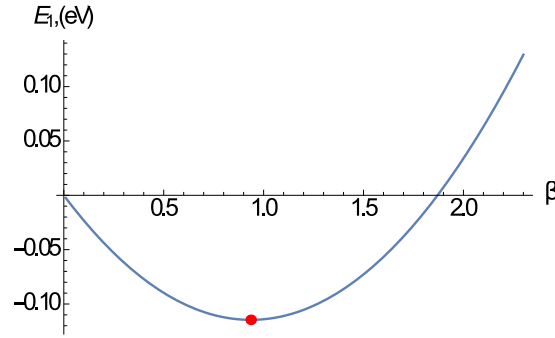


Figure 1. (Color online) The dependence of ground-state energy of extrinsic large polaron on the variational parameter β for $\eta = 0.08$ in La-based cuprates. The red point indicates a single carrier self-trapped state.

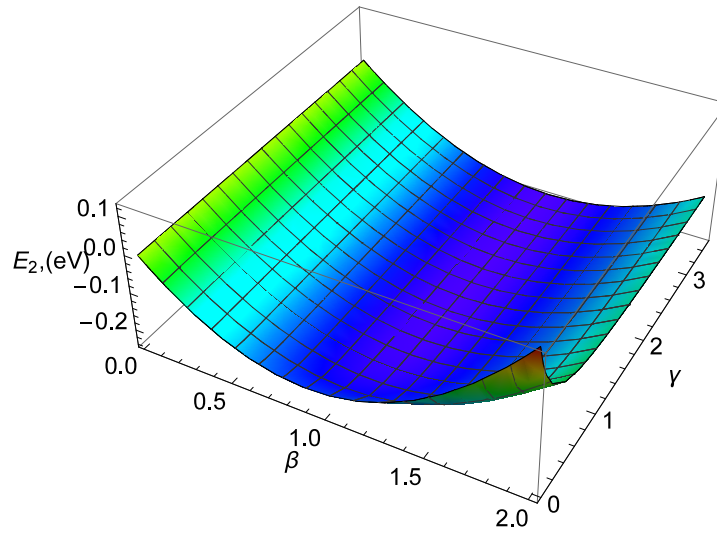


Figure 2. (Color online) The dependence of ground-state energy of extrinsic large bipolaron on the variational parameters β and γ for $\eta = 0.08$ in La-based cuprates.

Bohr radius $a_H = 0.529\epsilon_0(m_e/m^*) \text{ \AA}$ and the ionization energy $E_1^H = e^2/2\epsilon_0 a_H$ in lightly doped LSCO and LBCO. A similar argument can be made for other cuprates, such as $\text{La}_2\text{CuO}_{4+\delta}$, $\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ (YBCO) and $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ (Bi-2212). In these systems, the doping centers are excess oxygen atoms which capture the electrons from the oxygen conduction band and form the acceptor centers, i.e., O^- ions. One can assume that such extra O^- ions just like Sr^{2+} ions may have negative E_{dD} .

Figure 3 shows the variation of the ratio $E_{bU}/2E_{pI}$ with η for $b_s = 0.5$ and 2.0 for the stability region of the extrinsic large bipolaron in 3D cuprates. One can find out that such bipolarons exist as long as η is less than the critical value $\eta_c = 0.127$ and the ratio $E_{bU}/2E_{pI}$ reaches up to 0.287 (at $\epsilon_\infty = 3$ and $\eta \rightarrow 0$). We have determined the stability region of the extrinsic large bipolaron in cuprates and found that such bipolarons exist as long as η is less than the critical value $\eta_c = 0.127$ and the ratio $E_{bU}/2E_{pI}$ reaches up to 0.287 (at $\epsilon_\infty = 3$ and $\eta \rightarrow 0$). Figure 4 shows the variation of the ratio $E_{bB}/2E_p$ with η for $\epsilon_\infty = 3, 4$ and 5 for real large bipolarons in cuprates. One can see that in 3D cuprates, the intrinsic large bipolarons can exist at $\eta = \eta_c \leq 0.138$ and the ratio $E_{bB}/2E_p$ reaches up to 0.27 (at $\epsilon_\infty = 3$ and $\eta \rightarrow 0$). While the inset in figure 4 shows the variation of the ratio $E_{bB}/2E_p$ [calculated from equations (2.5) and (2.6) at $g_s = 0$ and $Z = 0$] with η for large optical bipolaron. Of course, besides the ratio $\eta = \epsilon_\infty/\epsilon_0$, the values of the Fröhlich electron-phonon coupling constant α are very important for the formation of large optical bipolarons. The long-range coupling of carriers with optical phonons is much stronger than their short-range coupling with acoustic phonons. Therefore, the long-range Fröhlich-type EPI in polar materials

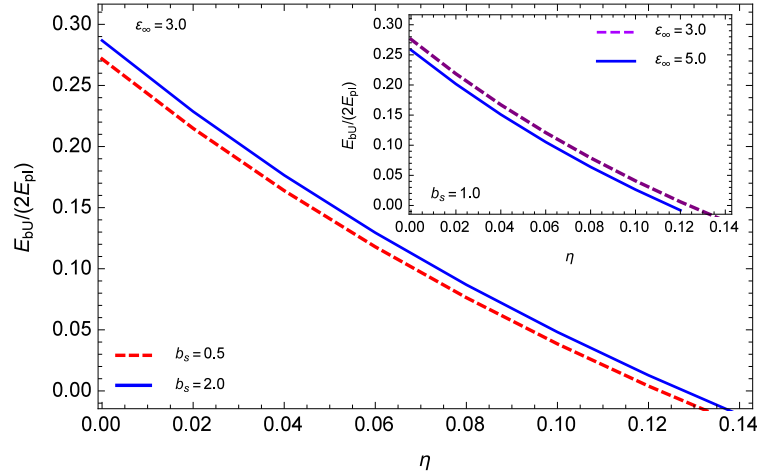


Figure 3. (Color online) The ratio of the binding energy of the extrinsic large bipolaron to twice the extrinsic large polaron binding energy as a function of η for two values of b_s in 3D cuprates. The inset shows the dependence of the ratio of the binding energy of the extrinsic large bipolaron to twice that of the extrinsic large polaron on η for two values of ϵ_∞ in 3D cuprates.

have been studied extensively [28, 29], although the short-range deformation potential type interaction is also important and leads to new effects. The dimensionless EPI Fröhlich coupling constant is defined as

$$\alpha = \frac{e^2(1-\eta)}{2\epsilon_\infty\hbar\omega_{LO}} \left(\frac{2m^*\omega_{LO}}{\hbar} \right)^{1/2}, \quad (2.7)$$

where ω_{LO} is the frequency of the longitudinal-optical (LO) phonon in an ionic crystal.

In polar materials, the formation of optical bipolarons is favored by larger values of α and by smaller values of η [2, 29], i.e., the optical (Fröhlich or Pekar) bipolarons exist only if α is greater than a critical value α_c and when $\eta < \eta_c$. Such 3D bipolarons can exist above rather high critical values α_c , e.g., $\alpha_c = 7.3$ as found by Adamowski [30] and $\alpha_c = 6.8$ found by Verbist, Peeters and Devreese [31]. Further, the value

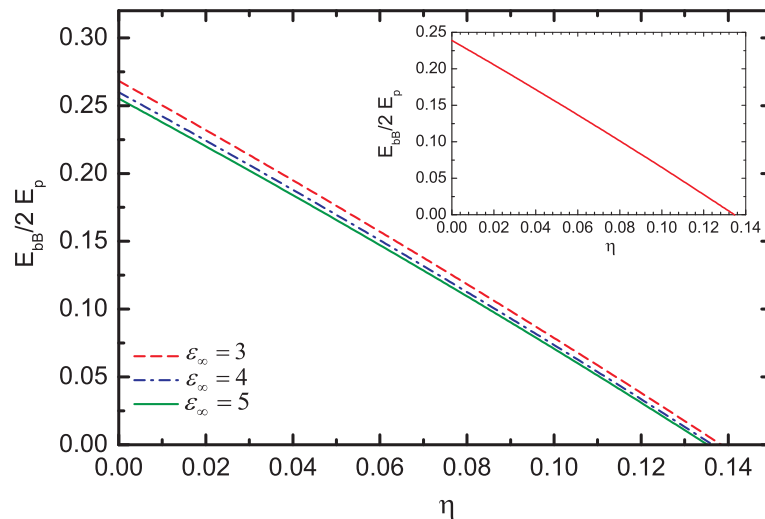


Figure 4. (Color online) The ratio of the binding energy of the real large bipolaron to twice that of the real large polaron as a function of η for different values of ϵ_∞ in 3D cuprates. The inset illustrates the ratio of the binding energy of the optical large bipolaron to twice that of the optical large polaron on η for $\epsilon_\infty = 3.0$ in 3D cuprates.

of α_c corresponding to the onset of the strong coupling regime is found to be $\alpha_c = 5.8$ [32], below which the formation of large optical bipolaron in 3D systems is unlikely. At a given value of η , the value of α_c depends on m^* , ϵ_∞ and ω_{LO} . The values of $\hbar\omega_{LO}$ in high- T_c cuprates range from 0.03 to 0.05 eV [13, 33]. Then, according to equation (2.7), the values of α corresponding to these high- T_c materials with $m^* = m_e$, $\epsilon_\infty = 3$ and $\eta = 0.02-0.10$ are equal to $\alpha = 4.96-6.95$. Thus, the conditions for the formation of large optical bipolarons are more favorable in the cuprates with $\epsilon_\infty = 3$, $\eta = 0.02-0.06$ and $\hbar\omega_{LO} = 0.03-0.04$ eV, at which conditions $\alpha > \alpha_c \approx 5.8$ [32] and $\alpha > \alpha_c = 6.8$ [31] are well satisfied.

We note here that the largest values of $\eta_c = 0.079-0.14$ and $E_{bB}/2E_p = 0.22-0.25$ as found in the literature [34–36] (see also [29]) were obtained for the optical bipolaron and do not correspond to specific substances. At the same time, the critical values of η below which the optical bipolarons can exist in cuprates were small enough ($\eta_c = 0.040-0.055$), as estimated in [29]. Our results are quite impressive in the sense that both the real and the optical bipolaron in 3D cuprates, can really exist for relatively large values of η (figure 4) and the large bipolarons are formed with the binding energies reaching up to 27% (at $\eta \rightarrow 0$) of twice the large polaron energy. The distinctive feature of the cuprates is their very large ratio of static to high-frequency dielectric constants. This situation is favorable for carriers attracted to polarization well created by the other ones or to Coulomb centers (dopants) to form 3D intrinsic or extrinsic large bipolarons. At $\epsilon_0 > 30$, such bipolarons (pair states) can be formed in lightly doped cuprates and they become unstable in an underdoped regime. As can be seen from tables 1 and 2, the binding energies of extrinsic and intrinsic large (bi)polarons and the ratios $E_{bU}/2E_{pI}$ and $E_{bB}/2E_p$ depend on several parameters. In particular, E_{pI} , E_p , E_{bU} and E_{bB} would rapidly increase with ϵ_∞ decreasing from 5 to 3, while the ratios $E_{bU}/2E_{pI}$ and $E_{bB}/2E_p$ increase more slowly as ϵ_∞ decreases. Interestingly, E_{pI} is an increasing function of η (table 1), while E_p is a decreasing function of η (table 2). Furthermore, we find that both E_{bU} and E_{bB} decrease with increasing η . We also find that the ratio $E_{bU}/2E_{pI}$ appreciably increases with b_s as shown in figure 3. The calculated values of the radii of extrinsic large (bi)polarons R_I (R_{BI}) (at $Z = 1$ and $b_s = 1$) and intrinsic large (bi)polarons, R_p (R_B) (at $Z = 0$ and $b_s = 0$) for different values of ϵ_∞ and η are presented in table 3.

Table 2. Calculated parameters of the real large polarons and bipolarons (with correlation between the pairing carriers) in 3D cuprates at different values of ϵ_∞ and η .

η	$\epsilon_\infty = 3.5$				$\epsilon_\infty = 4$				$\epsilon_\infty = 4.5$			
	E_p , eV	E_{bB} , eV	E_p^A , eV	E_B^A , eV	E_p , eV	E_{bB} , eV	E_p^A , eV	E_B^A , eV	E_p , eV	E_{bB} , eV	E_p^A , eV	E_B^A , eV
0.00	0.1107	0.0583	27.275	19.989	0.0845	0.0439	28.049	21.063	0.0666	0.0343	28.656	21.744
0.02	0.1063	0.0484	27.398	20.269	0.0811	0.0364	28.158	21.313	0.0640	0.0284	28.753	21.980
0.04	0.1019	0.0389	27.521	20.551	0.0778	0.0293	28.267	21.564	0.0614	0.0228	28.851	22.217
0.06	0.0977	0.0299	27.645	20.834	0.0746	0.0225	28.376	21.815	0.0588	0.0175	28.949	22.453
0.08	0.0935	0.0215	27.769	21.118	0.0714	0.0161	28.485	22.068	0.0563	0.0125	29.046	22.690
0.10	0.0895	0.0136	27.893	21.404	0.0683	0.0101	28.595	22.321	0.0539	0.0078	29.144	22.928
0.12	0.0855	0.0061	28.017	21.691	0.0653	0.0044	28.704	22.576	0.0515	0.0003	29.242	23.166
0.14	0.0816	—	28.142	21.980	0.0623	—	28.814	22.832	0.0491	—	29.341	23.404

Table 3. Calculated values of the radii of intrinsic and extrinsic large (bi)polarons R_p , R_I , R_B and R_{BI} in 3D cuprates at different values of ϵ_∞ and η .

η	$\epsilon_\infty = 3.5$				$\epsilon_\infty = 4$				$\epsilon_\infty = 4.5$			
	R_p , Å	R_B , Å	R_I , Å	R_{BI} , Å	R_p , Å	R_B , Å	R_I , Å	R_{BI} , Å	R_p , Å	R_B , Å	R_I , Å	R_{BI} , Å
0.00	8.6096	13.045	8.2839	12.403	9.8840	15.057	9.5633	14.434	11.158	17.066	10.841	16.457
0.02	8.7917	13.577	7.9061	12.084	10.092	15.663	9.1322	14.072	11.392	17.746	10.357	16.052
0.04	8.9813	14.146	7.5586	11.777	10.309	16.311	8.7357	13.724	11.636	18.475	9.9111	15.662
0.06	9.1791	14.756	7.2378	11.482	10.535	17.007	8.3699	13.390	11.890	19.256	9.5001	15.288
0.08	9.3854	15.413	6.9409	11.199	10.770	17.756	8.0313	13.068	12.155	20.098	9.1197	14.928
0.10	9.6009	16.121	6.6652	10.926	11.017	18.564	7.7169	12.758	12.432	21.005	8.7665	14.581
0.12	9.8261	16.887	6.4084	10.662	11.274	19.438	7.4243	12.460	12.722	21.986	8.4379	14.248
0.14	10.062	17.718	6.1688	10.408	11.543	20.386	7.1512	12.173	13.025	23.051	8.1312	13.926

3. Discussion

We now make some remarks about the characteristic (i.e., binding) energies of large bipolarons in the cuprates. The binding energies of extrinsic and intrinsic bipolarons strongly depend on ε_∞ , ε_0 and η . The values of ε_∞ and ε_0 observed in high- T_c cuprates are 3–5 [4, 37, 38] and 33–105 [4, 39], respectively, so that the values of η range from 0.028 to 0.15. Using the values of $\varepsilon_\infty = 3$ and $\eta \approx 0.03$, we find $E_{bU} \approx 0.07$ eV and $E_{bB} \approx 0.061$ eV. If we take other experimental values of $\varepsilon_\infty = 4$ and $\varepsilon_0 \approx 50$ for the cuprates (see [2, 29]), then we obtain $E_{bU} \approx 0.017$ eV and $E_{bB} \approx 0.016$ eV at $\eta = 0.08$. Further, at $\varepsilon_\infty = 5$ and $\eta = 0.08$ we find $E_{bU} \approx 0.0097$ eV and $E_{bB} \approx 0.0102$ eV. Thus, the extrinsic and intrinsic bipolarons can be experimentally found in high- T_c cuprates in the energy ranges ~ 0.01 – 0.07 eV and ~ 0.01 – 0.06 eV, respectively. The binding energies of large polarons and bipolarons are manifested in the excitation spectra of the hole-doped cuprates as the temperature-independent low-energy gaps or pseudogaps, which are different from the high-energy CT gaps ($\Delta_{CT} \approx 1.5$ – 2.0 eV [13]) of the cuprates.

It is of interest to compare our results with experimental data on localized in-gap states (or bands) and energy gaps (which are responsible for the existence of insulating phase and are precursors to the pseudogaps observed in the metallic state) in hole-doped cuprates. The above extrinsic and intrinsic (bi)polaronic states as well as hydrogenic impurity states emerge in the CT gap of the cuprates. In the experiments, these localized states are displayed as the in-gap states. One can see that the value of $E_{pI} \approx 0.13$ eV obtained at $\varepsilon_\infty = 4$ and $\eta = 0.1$ (table 1) is consistent with experimental data for lightly doped $\text{La}_2\text{CuO}_{4+\delta}$ [13]. The in-gap impurity band observed in this system at 0.13 eV might be associated with the extrinsic large polarons. While the values of $E_p \approx 0.096$ – 0.105 eV (table 2) obtained at $\varepsilon_\infty = 3.5$ and $\eta = 0.04$ – 0.06 agree reasonably well with the large pseudogap value ~ 0.1 eV observed in LSCO [40]. One of the important experimental observations is that in LSCO, the flatband [41], which is ~ 0.12 eV below the Fermi energy for $x = 0.05$, moves upwards monotonously with increasing x , but the flatband is lowered as x decreases and loses its intensity in the insulating phase. Apparently, the flatband observed by ARPES in the lightly doped LSCO ($x = 0.05$) is the energy band of large polarons since the effective mass of carriers obtained from the analysis of the ARPES spectra is about $2.1m_e$ [41]. The values of R_p (table 3) are also in good agreement with the experimental values of the radii of polarons which vary from 6 to 10 Å in cuprates [13].

4. Conclusions

We have studied the possible mechanisms of carrier localization in inhomogeneous hole-doped cuprates. The quantitative theory of the impurity-assisted and phonon-assisted single particle and pair ST of hole carriers in 3D lightly doped cuprates is developed within the continuum model and adiabatic approximation. The possible mechanisms for carrier localization lead to the formation of extrinsic large (bi)polaronic states, the hydrogenic impurity states (i.e., impurities with loosely bound free carriers or large polarons) and intrinsic large (bi)polaronic states in the CT gap of the cuprates. We have variationally calculated the binding energies and radii of the extrinsic and intrinsic large polarons and bipolarons, taking into account the short- and long-range parts of the carrier-defect-phonon and carrier-phonon interactions. We have determined the stability region of the extrinsic large bipolaron in cuprates and found that such bipolarons exist as long as η is less than the critical value $\eta_c = 0.127$ and the ratio $E_{bU}/2E_{pI}$ (where E_{pI} and E_{bU} are the binding energies of the extrinsic large polaron and bipolaron, respectively) reaches 0.287 (at $\varepsilon_\infty = 3$ and $\eta \rightarrow 0$). We have obtained the conditions for the real large bipolaron stability and estimated the values of $E_{bB}/2E_p$ in 3D cuprates (where E_p and E_{bB} are the binding energies of the intrinsic large polaron and bipolaron, respectively).

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Природа зовнішніх і внутрішніх самозахопних носіїв струму в слабкозлегованих купратних високотемпературних напівпровідниках

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Теоретично вивчено природу зовнішніх і внутрішніх самозахопних носіїв струму в купратах. Енергії зв'язування і радіуси зовнішніх і внутрішніх великих поляронів і біполяронів в купратах обчислено варіаційно, використовуючи неперервну модель та адіабатичне наближення. Ми показали, що зовнішні і внутрішні тривимірні великі біполярони існують в слабкозлегованих купратах, відповідно, при $\eta = \epsilon_{\infty}/\epsilon_0 < 0.127$ і $\eta < 0.138$ [де $\epsilon_{\infty}(\epsilon_0)$ — оптична (статична) діелектрична стала].

Ключові слова: полярон, біполярон, само захоплення, високотемпературні надпровідники