INSTABILITY OF CONNECTED OSCILLATOR SYSTEM UNDER ACTION OF MULTIPLICATIVE FLUCTUATIONS

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Studying of many oscillatory and wave processes is reduced to the studying of dynamics of connected oscillators. Parameters of these oscillators may be under influence fluctuation forces. Instability in this case develops. Such instability can be named fluctuation parametrical instability. In systems with one degree of freedom (one oscillator) this instability is investigated. Increment of it is proportional to intensity of fluctuations. In the present work it is shown, that in systems with the big number degree of freedom (two, three connected oscillators) may be realized conditions when such increment essentially grows. Possible consequence of such increase of increment may be fast diffusion fluctuations in systems with the big number degree of freedom.

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1. INTRODUCTION

Oscillator is model, which widely is used for studying of different oscillation processes. For example, many of oscillation in plasma physics may be considered as connected oscillators. This may perform by distinguishing temporal and spacing parts in solution of according physical problem. Space part is described by differential equation in partial derivatives. This allows to define space structure of electromagnetic fields. On other hand temporal processes dynamic is described by means temporal ordinary differential equations, which may be considered as describing the system of connected oscillators. consideration is possible complicated in electrodynamical structures filled by plasma. For example, cavity filled by plasma. Such system may be considered as two connected oscillators. One of them is electromagnetic, i.e. this is natural oscillation of cavity, and other is oscillation of plasma. This system was considered. It was supposed that plasma density was subjected to fluctuation action.

The dynamic as linear as nonlinear oscillators has enough studied. The different physical systems are subjected action of fluctuations. In oscillation ones they may be as additive as multiplicative. One of the important results is that multiplicative fluctuations may cause instability. It increment usually is proportional to fluctuation level. Ref. [1,2] are devoted different physical processes that may be considered as connected oscillators.

But it may be expected increment increasing of such instability in some cases. We investigated some of these ones. We showed that fluctuations instability increment might be proportional to root of some power of fluctuation level. It may be realized at some values of parameters, which describe concrete physical system. Values of parameters are defined.

First of all we considered two connected oscillators with multiplicative fluctuations. On other hand, we studied three connected oscillators, where each of them is connected with neighboring. Such system may be employed for modeling of continuos matter. We also made some conclusion when increment increase may be realized.

2. THE SYSTEM OF TWO CONNECTED LINEAR OSCILLATORS

Let us consider two connected linear oscillators, frequency one of them is subjected of random perturbation action. The differential equations describing this system are following:

$$\ddot{\xi}_{1} + \xi_{1} = -\mu \Omega^{2} \xi_{2} - \alpha \mu \Delta \xi_{2},
\ddot{\xi}_{2} + (\Omega^{2} + \Delta) \xi_{2} = -\mu \xi_{1},$$
(1)

where ξ_1, ξ_2 — generalized coordinates, describing these oscillators, μ — connection coefficient $(0 \le \mu \le 1), \Omega$ — dimensionless frequency of the second oscillator (the dimensionless time is normalized on the first oscillator period, $\Delta << \Omega^2$ — random addition to second oscillator frequency, α — coefficient accounting additive influence of second oscillator on first one $(0 \le \alpha \le 1)$. Random addition is Gaussian and delta-correlation,

$$\langle \Delta (\tau) \Delta (\tau') \rangle = N \delta (\tau - \tau'), \qquad (2)$$

where N – noise level.

The differential equation set for first moments which are average values of $\xi_1, \dot{\xi}_1, \xi_2, \dot{\xi}_2$, coincides with set (1) when Δ =0, therefore they are not interest for us. We investigated second moments, which are average values of all products above mentioned $\xi_1, \dot{\xi}_1, \dot{\xi}_2, \dot{\xi}_2$. The amount of second moments is ten. The differential equations set for them was obtained. The technique of variational derivatives and Furutsu-Novikov formula was used for obtaining of second moments equations set. This technique is described in [2].

For solution defining of these equations in exponential form as $\exp(\lambda \tau)$, the characteristic equation was obtained. It is algebraic equation of tenth power. It may be presented such as consisting of two parts:

$$\prod_{i=1}^{10} (\lambda - \lambda^{(2)}) = NP(\lambda),$$

where $\lambda^{(2)}$ – roots of left part, N – fluctuation level. Left part is polynomial. Right part is polynomial too proportional to fluctuation level. The right part of characteristic equation in such form is conditioned by fluctuations influence. The roots of left part polynomial simply are connected with natural frequencies of initial connected linear oscillators system without fluctuations, which describing by set (1) when Δ =0. These are all different sums and differences of natural frequencies multiplied on imaginary one. The characteristic equation for (1) when N=0 is following:

$$\lambda^4 + (1 + \Omega^2)\lambda^2 + \Omega^2(1 - \mu^2) = 0$$
,

and its solutions are:

$$\lambda_{1,2}^{(0)} = \pm \frac{i}{\sqrt{2}} \sqrt{1 + \Omega^2 - \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}},$$

$$\lambda_{3,4}^{(0)} = \pm \frac{i}{\sqrt{2}} \sqrt{1 + \Omega^2 + \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}}.$$

They coincide with natural frequencies without imaginary one *i*. Accordingly for $\lambda^{(2)}$ we obtain following expressions:

$$\begin{split} \lambda_{1,2}^{(2)} &= \pm \sqrt{2} i \sqrt{1 + \Omega^2 - \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} \\ \lambda_{3,4}^{(2)} &= \pm \sqrt{2} i \sqrt{1 + \Omega^2 + \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} \\ \lambda_{5,6}^{(2)} &= 0 \\ \lambda_{7,8}^{(2)} &= \pm \frac{i}{\sqrt{2}} \left\{ \sqrt{1 + \Omega^2 - \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} + \right. \\ &\left. + \sqrt{1 + \Omega^2 + \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} \right. \\ \lambda_{9,10}^{(2)} &= \pm \frac{i}{\sqrt{2}} \left\{ \sqrt{1 + \Omega^2 - \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} - \right. \\ &\left. - \sqrt{1 + \Omega^2 + \sqrt{\left(1 - \Omega^2\right)^2 + 4\mu^2 \Omega^2}} \right. \end{split}$$

When noise is low, the roots of characteristic equation may be presented in following form: $\lambda = \lambda^{(2)} + \delta$. The addend δ is conditioned by fluctuation influence. λ is complex value. If real part of λ (i.e. real part δ) is positive, there is instability on the second moments in the system consisting of two connected linear oscillators. This instability is conditioned by fluctuations, and it increment equals real part. In general case this increment is proportional to noise level. But it is possible that δ will be proportional to root of some power of fluctuations level. It may be when there are multiple roots among ones of left part polynomial. In this case the addend δ is proportional to root of power, which coincide with the left part equal roots number. In general case the equal roots of left part of characteristic equation may exist at some parameters of initial set (1). The values of these parameters were defined. When left part of characteristic equation has multiple roots, increment of mentioned instability becomes more large comparatively the ordinary case. We defined increments of second moment instability, which were proportional as cubic root as square root of fluctuation level. The parameters values, when this is possible, were defined.

The above considered may be using for concrete physical system. It is cylindrical cavity filled by plasma with fluctuating density. Such system may be considered as two connected linear oscillators. One of them corresponds to electromagnetic oscillations of cavity other corresponds to plasma oscillations. The last is objected to plasma density fluctuations. There is instability on second moment in such system. But there is not expected increase of increment in cavity with fluctuating plasma.

3. SOME CONCLUSION ABOUT ARBITRARY OSCILLATORS NUMBER

Let us consider system of arbitrary number linear connected oscillators described by following differential equations set:

$$\ddot{x}_i + \omega_i^2 x_i = \sum_{j=1, j \neq i} \alpha_{ij} x_j \quad i = 1, 2, ..., n,$$
(3)

where x_i – generalized coordinate describing oscillator, ω_i - it frequency, a_{ij} - connection coefficients, n – oscillators number. We suppose that there is frequency fluctuation in one of oscillators. The above mentioned technique of variational derivatives may be used for moments equations set obtaining. It may be shown that first moments equation set coincides with set (3) without fluctuations. The second moments are interesting. The first moments number is $M_1 = 2n$. It may be shown that second moments number equals $M_2 = n(2n+1)$. The differential equations set may be obtained. It will be linear homogeneous set. Solutions may be presented in exponential form as for two oscillators. There is characteristic equation for second moments. It is algebraic equations of M_2 power. It may be presented in form of two parts, as that performed for two oscillators. The left part of this equation may have more large number of multiple roots than in two oscillators case. The maximum multiplicity is defined by as oscillators number as parameters number describing this physical system. Maximum multiplicity may equal $k = \min(M_2, M_p)$, where M_p – number of parameters of considered system.

In general case the system containing large number of connected linear oscillators and described by large number of parameters, and subjected action of multiplicative fluctuations may have large multiple roots number of characteristic equation left part for second moments. Solutions of characteristic equation in the case of low fluctuation level may be presented in form $\lambda = \lambda^{(2)} + \delta$, as it was mentioned above. In the case of left part roots with large multiplicity the addend δ will be proportional to root of multiplicity power of fluctuations level. The addend δ may be larger, if multiplicity of corresponding root is large, in the case of many oscillators comparatively case of two oscillators. It was mentioned that the roots λ are complex values. Therefore δ is complex too. If real part δ is positive there is instability on second moments in system of many oscillators with multiplicative fluctuations. If multiplicity of mentioned root is large the instability increment, which equals real part of δ , will be large too. In the case of many oscillators it may be more larger comparatively case of two connected linear oscillators.

4. THREE CONNECTED OSCILLATORS UNDER ACTION OF MULTIPLICATIVE FLUCTUATIONS

In this section we shell be shown that in system of three connected oscillators with multiplicative fluctuations the essential increase of fluctuation instability is possible. Such system is describing by following:

$$\ddot{x}_{1} + \omega_{1}^{2} x_{1} = \alpha_{12} x_{2} ,$$

$$\ddot{x}_{2} + \omega_{2}^{2} x_{2} = \alpha_{21} x_{1} + \alpha_{23} x_{3} ,$$

$$\ddot{x}_{3} + \omega_{3}^{2} x_{3} = \alpha_{32} x_{2} ,$$
(4)

where symbols are as in (3). The characteristic numbers of this system are:

$$\lambda_{1,2} = \pm i, \ \lambda_{3,4} = \pm i\sqrt{1+A}, \ \lambda_{5,6} = \pm i\sqrt{1-A},$$

where $A = a_{12}a_{21} + a_{23}a_{32}$.

We suppose that frequency of first oscillators ω_1 is under action of fluctuations, which are delta-correlative with level N_1 . The number of second moments is 21. The differential equations set for them was obtained. The technique of variational derivatives and Furutsu-Novikov formula as in section 2 was used for obtaining of second moments equations set too. As in section 2, the characteristic equation was obtained. It was presented in form:

$$\prod_{i=1}^{21} (\lambda - \lambda^{(2)}) = N_1 P(\lambda). \tag{5}$$

 $\lambda^{(2)}$ are all possible sums of λ_i . The roots were analyzed. The goal of this analysis was to obtain the existence multiple roots among $\, \lambda^{\, (2)} \, . \,$ There are roots with different multiplicity. The parameters values when this is possible were obtained. $\lambda^{(2)}=0$ with maximal multiplicity that equals 9 and $\lambda^{(2)} = \pm 2i$ with multiplicity 6. It possible when parameter A=0. But such $\lambda^{(2)}$ may be roots of polynomial $P(\lambda)$. To exhibit this the right part of (5) was transformed. This right part is determinant. The determinant order was reduced to order of three. It was exhibited that polynomial $P(\lambda)$ has multiple roots coinciding with $\lambda^{(2)}$. Just it is 0. As mentioned above in section 2 the addend δ is proportional to root of multiplicity power of noise level for multiple $\lambda^{(2)}$. But this condition is violated when polynomial $P(\lambda)$ has same roots. In this case the root index in the expression for δ is less on number that equals multiplicity of corresponding root of $P(\lambda)$. In our case such root is $\lambda^{(2)}$ =0. λ =0 is fourth degenerated root of polynomial $P(\lambda)$.

Therefore $\delta \approx \sqrt[5]{N_1}$ for $\lambda^{(2)} = 0$ and $\delta \approx \sqrt[6]{N_1}$ for $\lambda^{(2)} = \pm 2i$. There are complex δ with positive real part. There is instability in this case. Fluctuations are faint. Therefore we have fluctuation increment increasing when multiple roots $\lambda^{(2)}$ exist. Thus as it see from this section there is increasing possibility of fluctuation instability increment for three connected oscillators comparatively of two ones case. This confirms conclusions made in section 3. Results of this section may be used when fluctuation influence on continuos matter is investigated. Continuos matter may be presented as chain of connected oscillators.

5. SUMMARY

It is well known that multiplicative fluctuations may cause instability. It was shown on the example of one oscillator, which frequency is subjected to random action, in [2]. This instability exists on second moments for deltacorrelation fluctuation. In this repot possibility of fluctuation instability increment increasing considered. It has been noted that such increasing may take place in a system of large oscillators number, which described by large parameters number. It was considered on the examples of two and three connected oscillators with multiplicative fluctuations. It is necessary to note that in the case of three oscillators instability increments may be more large. The necessary conditions of increment increasing were investigated. The parameters when increment increasing takes place were defined.

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НЕСТІЙКІСТЬ ЗВ'ЯЗАНИХ ОСЦИЛЯТОРІВ, ЯКІ ЗНАХОДЯТЬСЯ ПІД ДІЄЮ МУЛЬТІПЛІКАТИВНИХ ФЛУКТУАЦІЙ

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Вивчення багатьох коливальних та хвильових процесів приводить до вивчення динаміки зв'язаних осциляторів. Параметри цих осциляторів можуть опинитися під дією флуктуаційних сил. В цьому випадку розвивається нестійкість. Таку нестійкість можна назвати флуктуаційною параметричною нестійкістю. В системах з одним ступенем свободи (один осцилятор) ця нестійкість вивчена. ЇЇ інкремент пропорційний інтенсивності флуктуацій. В даній роботі показано, що в системах з великим числом ступенів свободи (два, три зв'язаних осциляторів) можуть бути реалізовані умови, коли такий інкремент суттєво зростає. Можливим слідством такого зростання інкременту може бути швидке розсасування флуктуацій в системах з великим числом ступенів свободи.

НЕУСТОЙЧИВОСТЬ СВЯЗАННЫХ ОСЦИЛЛЯТОРОВ, НАХОДЯЩИХСЯ ПОД ДЕЙСТВИЕМ МУЛЬТИПЛИКАТИВНЫХ ФЛУКТУАЦИЙ

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Изучение многих колебательных и волновых процессов сводится к изучению динамики связанных осцилляторов. Параметры этих осцилляторов могут оказаться под воздействием флуктуационных сил. В этом случае развивается неустойчивость. Такую неустойчивость можно назвать флуктуационной параметрической неустойчивостью. В системах с одной степенью свободы (один осциллятор) эта неустойчивость изучена. Инкремент ее пропорционален интенсивности флуктуаций. В настоящей работе показано, что в системах с большим числом степеней свободы (два, три связанных осциллятора) могут быть реализованы условия, когда такой инкремент существенно возрастает. Возможным следствием такого возрастания инкремента может быть быстрое рассасывание флуктуаций в системах с большим числом степеней свободы.