

HYSTERESIS EFFECTS IN ELECTRON GUNS

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Computer simulations of beam formation in an electron gun with large compression of a thin annular beam show some hysteresis effects: a beam with different currents can exist in the same external conditions. The results of computer simulation are presented for several simple drift and diode geometry and for the gun with large compression of an annular beam. Physics reasons of such behaviour are discussed.

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1. INTRODUCTION

Recently one begun to use photocathodes that permit forming "cold" beams in sources of electron beams of high brightness developed and applied for FELs and linear colliders. One of the important problems in these investigations concerns the value of the current taken from the photocathode. In a number of works – theoretical as well as experimental – there is discussed the possibility of obtaining currents exceeding the Child-Langmuir limit and is indicated that such an increase is possible only for the generation of short bunches the length of which is less or of the order of the gap length.

Naturally, the classical CL-law is applicable only for a stationary regime with the current limited by the space charge for the following conditions: 1) voltage constant, 2) transverse dimensions of the emitting surface significantly exceeding (geometry close to unidimensional) the characteristic longitudinal dimensions of the accelerating region and 3) emissive power of the cathode unlimited.

Let us consider two first of the limitations. The first corresponds to the current constancy and implicitly implies that the variable quantity can be only the voltage, while the current value determined by this voltage. It should be noted that with the appearance of efficient photoemitters and lasers, it became possible to vary independently the emission current, in addition to the possibility of varying the voltage. The possibility of controlling the beam current in photoemission sources at a fixed voltage leads to a qualitatively different situation. To a certain extent it is suggestive of the process of beam injection (although at a very low initial energy) in an accelerating region rather than drawing the beam from an emitting surface.

The second limitation reduced the problem to unidimensional and, evidently, is violated under real conditions. Although the influence of the end effects is recognised, its contribution is considered to be small. At the same time, in a plane diode, when the transverse dimensions of the emitter are approximately equal to the length of the gap, the influence of fringe effects becomes determining. They are no longer simply fringe effects, they are main effects.

In view of the non-stationary nature and multidimensionality of the problem, the results given below were obtained by computer simulation with electromagnetic code KARAT [1].

2. FRINGE EFFECTS IN SIMPLE DRIFT AND ACCELERATING REGIONS

For a purely stationary flow that is one-dimensional and unlimited in the transverse direction, the dependence of the passing current I_{out} on injection current I_{inj} is represented by the so-called "lambda-curve" (Fig. 1) [2].

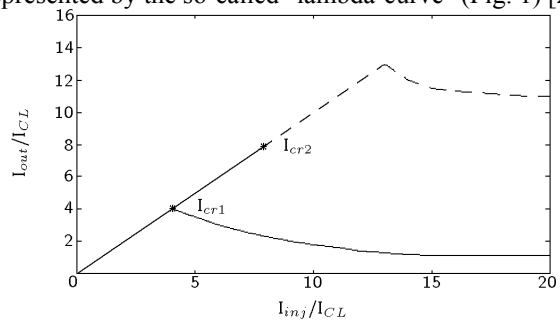


Fig. 1. λ -curve

Three solutions are possible for the region from I_{cr1} to I_{cr2} in such a flow. One of these is with a virtual cathode and partial passage of the beam (lower part of solid curve) and the other two (merging into one) are without a virtual cathode and with complete passage of the beam. It is considered that the last two states are unstable and when I_{cr1} is exceeded the beam always goes to the lower part of the curve with formation of a virtual cathode. For unlimited increase of injected current the virtual cathode asymptotically approaches the plane of injection and through the region there passes only current equivalent to the CL-current for a plane diode of the same dimensions, with voltage corresponding to the energy of injected electrons. For non-relativistic energy of particles, $I_{cr1} = 4 I_{CL}$, $I_{cr2} = 2 I_{cr1}$.

In [3], it was analytically shown that when the flow completely fills the chamber cross-section (the case of a semi-infinite plane slit with end plane was considered), the current passed can increase without limit $I_{out} \propto (I_{inj})^{1/2}$ due to "exudation" of particles close to the chamber wall in a thin layer. The thickness of this layer decreases inversely proportionally to $(I_{inj})^{1/2}$. Close to the end plane, the central part of the beam is blocked by the space-charge field and the virtual cathode acquires the form of such a coaxial slit with rounded edges.

Let us consider the case when two-dimensional effects appear most effective. The drift region is chosen in the form of a cylindrical resonator having the diameter ($D=5$ cm) much larger than the longitudinal dimension ($d=1$ cm), the diameter of injected beam ($2r_e=1$ cm)

equals to the length of the resonator. We present the results for a beam having non-relativistic energy $W=20$ keV. In Fig.1, the value of current is presented in units corresponding to CL-current ($I_{CL} = \pi r_e^2 j_{CL}, j_{CL}d^2 = 6.6$ A for 20 kV on the diode). The solid line show the analytical dependence for one-dimensional stationary flow and the dotted lines the calculated dependence for established regime when injecting beams of various currents with a 4-ns linear front. The established beam configuration for 200-A injected current is shown in Fig.2 and is very similar to that considered in [3].

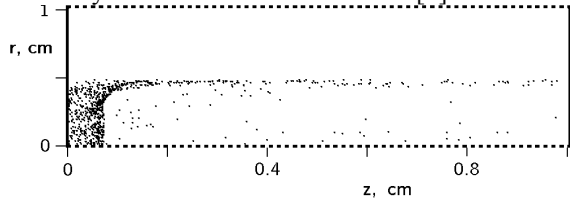


Fig.2. Beam with "whiskers"

The difference is that the beam passes far from the side walls and the passed current does not increase with the injected, but remains approximately at one level. It should be noted that the main part of the current is concentrated in "whiskers", the thickness of which decreases with increasing injected current.

The time dependence of the current passing through the resonator is shown in Fig.3. The interruption of current is due to the formation of a virtual cathode arising near the center of the resonator and is rapidly displaced toward the injection plane when reflected electrons appear.

With increasing transverse dimensions the influence of "whiskers" on the inner part of the beam is weakened due to preventing of field by electrodes and electrons with density of current close to the Child-Langmuir begin take "emitted" from the virtual-cathode region. Since the total CL-current will increase proportionally to r^2 and the current in the "whiskers" not more than r , the contribution of the "whiskers" becomes small relative to current from a larger area, but lower current density. This corresponds to one-dimensional flow unlimited in the transverse direction.

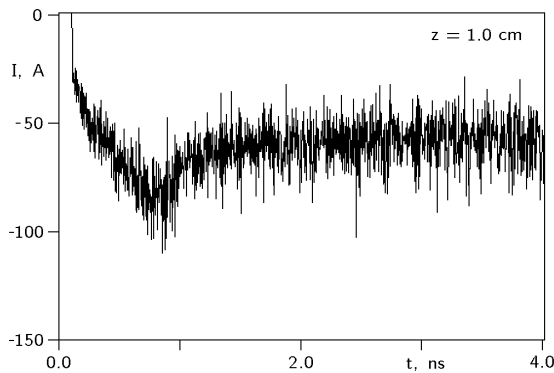


Fig.3. Transient regime in the drift region

The detailed exposition of the effects occurring in a drift region had to show, as its main purpose, that qualitative concepts about the formation of beam structure in the drift region can be transferred to the accelerating region. In any case, this is so when it is possible to rapidly control emission current, i.e., for the case of photocathodes.

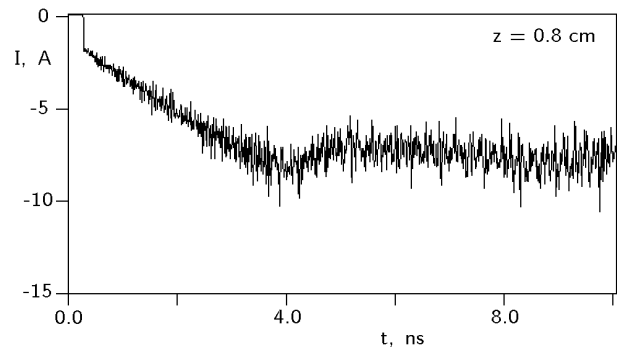


Fig.4. Transient regime in the plane diode

The formation of "whiskers" in a drift region occurs particularly effectively when near the cathode boundary there is a conducting wall and particles simply drift along it with the energy close to the energy of injection. In the case of an accelerating region, the energy of emitted electrons is low and the formation of fringe "whiskers" becomes difficult because of the low energy of "injection" and strong influence of the proper field. The proximity of the electrode, at the cathode potential, directly to the boundary of the emitting surface leads to screening of the proper as well as extracting field. Thus, it is necessary that the distance of side walls from emitter edges be such as to allow, on the one hand, considerable screening the influence of own space charge and, on the other, penetration of a sufficiently large external field.

As an example that supports such an approach, we present in Fig.4,5 time dependencies of the beam current in a plane diode and in a diode with emitting surface in a recess. The voltage on the diode is constant and equal to 20 kV, emitter diameter is 1 cm, emission current increases linearly to 8 A in 4 ns and is then held constant. In the first case, the hysteresis is practically imperceptible, but in the second case it is quite pronounced.

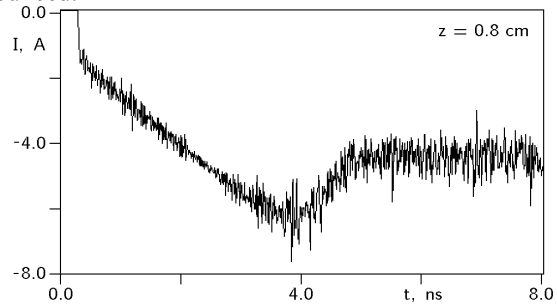


Fig.5. Transient regime in the conical diode

3. HYSTERESIS EFFECTS IN THE HIGH-VOLTAGE ELECTRON GUN

Spatial-temporal hysteresis effects to a significant extent appeared in the process of developing an electron gun with large compression of a tubular beam (Fig.6) intended for generation of powerful microwave and IR radiation [4].

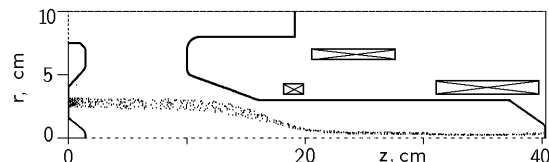


Fig. 6. Schematic diagram of the gun

An electron gun has been designed to produce a beam for generation of high-power microwaves and IR-radiation. It has the following design parameters: energy of electrons 350...450 keV, beam current >100 A, longitudinal momentum spread <1%, pulse duration 10 μ s, repetition rate up to 10 Hz.

A plane, ring-shaped, cathode, having a 28-mm average radius and a 6...8 mm width, is used to produce an annular beam. The cathode surface is screened by two focusing electrodes to operate in a space-charge limited regime. The focusing system must form a laminar beam inside the accelerating gap at rather low magnetic field and transport the beam to an experimental region of about 50-cm length. Within this region, the beam radius is 2.5 mm and annular thickness is less than 1 mm due to a large magnetic field of 15 kGs.

The code KARAT was used to calculate the beam dynamics and the results were tested by means of the stationary code SAM [5]. Simulation of emission of particles in the KARAT code could be realised by two methods: giving constant (or variable) emission current exceeding the space-charge limited current and increasing the voltage on the gun to a certain given value (thermionic cathode); giving constant voltage on the gun and increasing emission current from the cathode in time in accordance with a desired law (photocathode).

Fig. 7 shows the averaged (high-frequency components filtered off) behaviour of beam current at the gun output for various modelling conditions. Dependencies 1-3 correspond to conditions of photocathode operation and dependence 4 to conditions of thermionic cathode operation. Dependencies 1-3 were obtained for the constant rate of emission current increase (about 28 A/ns) on the linear front, the duration of which was 7 ns (195 A), 7.5 ns (210 A) and 11 ns (310 A) for dependencies 1-3, respectively, and the constant current for greater time. Current discrimination is clearly seen at 195 A, with subsequent drop to values that approximately correspond to the beam current value in the stationary state obtained with stationary code SAM. The fourth dependence is plotted for the case when the emis-

sion current is held constant at 400 A while the voltage on the diode increases from 160 kV to 400 kV in 8 ns.

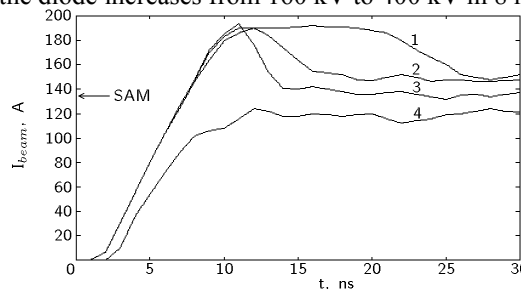


Fig. 7. Hysteresis effects in the high-voltage gun

It should be noted that a little screening of the cathode (for example, placing it in a 5-mm deep circular slot), leading of course to a decrease in diode perveance, is accompanied by a sharper change in the ratio of peak to stationary value of current.

4. CONCLUSIONS

Computer simulations of beam formation in the gun show some hysteresis effects: a beam with a different current can exist in the same external conditions. At the upper (metastable) part of a hysteresis curve the beam current exceeds the Child-Langmuir current. Experimental observation of these metastable states is possible if one has the possibility for fast control of emitted current, which is the case for photoemission cathode.

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ГИСТЕРЕЗИСНЫЕ ЯВЛЕНИЯ В ЭЛЕКТРОННЫХ ПУШКАХ

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При исследовании нестационарного режима формирования тонкостенного электронного пучка с большой компрессией по радиусу обнаружены гистерезисные эффекты: при одних и тех же внешних условиях ток пучка мог принимать разные значения. Приведены результаты численного моделирования формирования пучков для нескольких простых дрейфовых и диодных геометрий, и для пушки с большой компрессией пучка. Обсуждаются физические причины гистерезисных эффектов. Работа выполнена при поддержке РФФИ по гранту 03-02-17301.

ГИСТЕРЕЗИСНІ ЯВИЩА В ЕЛЕКТРОННИХ ПУШКАХ

А.В. Агафонов, А.М. Лебедев

При дослідженні нестационарного режиму формування тонкостінного електронного пучка з великою компресією по радіусі виявлені гистерезисні ефекти: при тих самих зовнішніх умовах струм пучка може приймати різні значення. Приведено результати чисельного моделювання формування пучків для декількох простих дрейфових і діодних геометрій, і для гармати з великою компресією пучка. Обговорюються фізичні причини гистерезисних ефектів. Робота виконана за підтримкою РФФД по гранту 03-02-17301.