

Anisotropic g factors of the tetragonal Cu^{2+} monomer in TI-2223 superconductor

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The gyromagnetic factors of the Cu^{2+} monomer in TI-2223 superconductor are quantitatively investigated from the perturbation formulas of these factors for a $3d^9$ ion in a tetragonally elongated octahedron. The local tetragonal distortion of the system is attributable to the axial elongation along c axis, corresponding to the five-fold coordinated $\text{Cu}^{2+}(2)$ site with almost 30% longer Cu–O bond length for the apical oxygen as compared to the four planar ones. The significant anisotropic behaviors of the EPR spectra perpendicular to and parallel with the ab (CuO_2) layers are analyzed on the basis of the local tetragonal elongation.

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Introduction

$\text{Tl}_2\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_{10}$ (TI-2223) superconductor has been extensively investigated due to the interesting electrochemical [1], specific heat [2], dielectric [3], microwave [4] and superconductive [5–7] properties. It is well known that these properties may be related to the local structures and electronic states of the CuO_2 planes in TI-2223 and can be studied by means of the electron paramagnetic resonance (EPR) technique. The EPR experiments were carried out for the Cu^{2+} monomer in TI-2223 after air sucking at 100°C for 12 h, with anisotropic g factors g_{\parallel} and g_{\perp} measured at room temperature [8]. Until now, however, the experimental results have not been theoretically explained.

Since the EPR behaviors and microscopic mechanism of the Cu^{2+} monomer in TI-2223 would be useful to understand the magnetic and superconductive properties of TI-2223, further investigations on the EPR g factors are of fundamental and practical significance. On the other hand, the local structure of Cu^{2+} site can be important for the analysis of the EPR behaviors of EPR. Furthermore, the changes of bond lengths bring forward the obvious influences on the crystal field and hence g factor. In this work, the EPR spectra of the Cu^{2+} monomer in TI-2223 are quantitatively studied using the high order perturbation formulas of the g factors for a tetragonally elongated octahedral $3d^9$ -cluster. In the calculations, the local structure of the five-fold coordinated $\text{Cu}^{2+}(2)$ site is correlated to the EPR analysis from the superposition model.

Calculations

The observed EPR signal may be ascribed to the tetragonal five-fold coordinated $\text{Cu}^{2+}(2)$ site in TI-2223, with one longer axial Cu–O distance R_{\parallel} ($\approx 2.48 \text{ \AA}$) and four shorter planar bond lengths R_{\perp} ($\approx 1.927 \text{ \AA}$ [9], the structural data at 295 K). This $[\text{CuO}_5]^{8-}$ complex can be described as a square pyramid (or half-octahedron without one vertex), i.e., one of original oxygen in an ideal octahedron is stretched to infinity and thus leads to an extremely elongated octahedron (see Fig. 1 of Ref. 10). For a $3d^9(\text{Cu}^{2+})$ ion in a tetragonally elongated octahedron, its original cubic ground 2E_g level would be separated into two orbital singlets ${}^2A_{1g}$ and ${}^2B_{1g}$, with the latter lying lowest, while the cubic excited ${}^2T_{2g}$ state may split into an orbital singlet ${}^2B_{2g}$ and a doublet 2E_g [11]. The high order perturbation formulas of the g factors for a $3d^9$ under tetragonally elongated octahedra can be expressed as [12]:

$$\begin{aligned} g_{\parallel} &= g_s + 8k\zeta/E_1 + k\zeta^2/E_2^2 + 4k\zeta^2/(E_1E_2) + \\ &+ g_s\zeta^2[1/E_1^2 - 1/(2E_2^2)] + k\zeta^3[4/(E_1E_2^2) - 1/E_2^3] - \\ &- 2k\zeta^3[2/(E_1^2E_2) - 1/(E_1E_2^2)], \\ g_{\perp} &= g_s + 2k\zeta/E_2 + k\zeta^2[2/(E_1E_2) - 1/E_2^2 - 4/(E_1E_2)] + \\ &+ 2g_s\zeta^2/E_1^2 + k\zeta^3(4/E_1^2 - 1/E_2^2)/(2E_2), \end{aligned} \quad (1)$$

where g_s (≈ 2.0023) is the pure spin value. k and ζ are, respectively, the orbital reduction factor and the spin-orbit coupling coefficient for the $3d^9$ ion in crystals. The denominators E_1 and E_2 denote the energy separations between the excited ${}^2B_{2g}$ and 2E_g and the ground ${}^2B_{1g}$ states [12], which can be determined from the energy matrix of a $3d^9$ ion under tetragonal symmetry:

$$E_1 = 10Dq, \text{ and } E_2 = 10Dq - 3Ds + 5Dt. \quad (2)$$

Here Dq is the cubic field parameter, and Ds and Dt are the tetragonal ones. The tetragonal field parameters may be calculated from the superposition model [12] and the local geometry of this $\text{Cu}^{2+}(2)$ site in TI-2223:

$$Ds \approx (4/7)\bar{A}_2 [(R/R_\perp)^{t_2} - (1/2) (R/R_\parallel)^{t_2}],$$

$$Dt \approx (16/21)\bar{A}_4 [(R/R_\perp)^{t_4} - (1/2) (R/R_\parallel)^{t_4}]. \quad (3)$$

Here t_2 (≈ 3) and t_4 (≈ 5) are the power-law exponents [13]. \bar{A}_2 and \bar{A}_4 are the intrinsic parameters, with the reference distance R taken as the average of the parallel and perpendicular Cu–O distance. For $3d^9$ ions in octahedra, the relationships $\bar{A}_4 \approx (3/4) Dq$ and $\bar{A}_2 \approx 10.8 \bar{A}_4$ [13] have been proved valid in many crystals and are suitably employed here. So the EPR g factors, especially the anisotropy Δg ($= g_\parallel - g_\perp$), are correlated to the tetragonal field parameters and hence to the local structure of the $\text{Cu}^{2+}(2)$ site.

From the optical absorption data for Cu^{2+} in oxides [14], the spectral parameters $Dq \approx 1260 \text{ cm}^{-1}$ and $k \approx 0.78$ can be obtained for the $\text{Cu}^{2+}(2)$ site in TI-2223. According to the free-ion value $\zeta_0 \approx 829 \text{ cm}^{-1}$ [11] for Cu^{2+} , the spin-orbit coupling coefficient $\zeta \approx k\zeta_0$ can be calculated. Substituting the values mentioned above into Eq. (1), the theoretical g factors as well as the anisotropy Δg are obtained and shown in Table 1.

Table 1. The g factors and anisotropy of the Cu^{2+} monomer in TI-2223

	g_\parallel	g_\perp	Δg
Calculation	2.350	2.050	0.300
Experiment [8]	2.35	2.10	0.25

Discussion

One can find from Table 1 that theoretical g factors for the Cu^{2+} monomer in TI-2223 based on the perturbation formulas Eq. (1) and the local structural parameters of the Cu^{2+} site show reasonable agreement with the observed values. Thus, the EPR experimental data [8] are satisfactorily interpreted for the first time in this work.

The low symmetrical (tetragonal) distortion of the system arises from the tetragonal elongation of this five-fold coordinated Cu^{2+} site in TI-2223, i.e., one longer axial R_\parallel ($\approx 2.48 \text{ \AA}$) and four shorter planar R_\perp ($\approx 1.927 \text{ \AA}$) [9] with about 30% relative discrepancy. Although the tetrag-

onal elongation is ascribed to the local structure of the Cu^{2+} site in TI-2223, the influence of the Jahn–Teller effect cannot be excluded. This point was mentioned for the similar Cu^{2+} center on the octahedral Ca^{2+} site in $\text{Ca}(\text{OH})_2$ [15].

The positive observed anisotropy Δg may be interpreted as the tetragonal elongation of the $[\text{CuO}_5]^{8-}$ complex, which results in the ground ${}^2B_{1g}$ state. Further, the significant anisotropic behaviors of the EPR spectra perpendicular to and parallel with the ab (CuO_2) layers are also understandable in view of the local tetragonal elongation. The theoretical studies in this work seem helpful to understand material microstructures and magnetic properties of CuO_2 plane in high T_c superconductors.

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