ON THE MECHANISM OF ${}^{4}He(\gamma, pn)^{2}H$ REACTION

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The interpretation is given on the effect of two "peaks" appearance in the experimental deuteron energy distributions in the ${}^{4}He(\gamma, pn)^{2}H$ reaction with photon energies $E_{\gamma}=30...50$ MeV. The mechanism of pole ${}^{2}H$ -diagram determining the first "peak" motion and the mechanism of triangular ${}^{3}He({}^{3}H)$ -diagrams showing a fixed appearance of the second "peak" with a deuteron energy $E_{2H}^{m}=2.22$ MeV due to the root singularity in the $p + n \rightarrow {}^{2}H$ vertex are taken into account. The position of maxima in the calculated distributions by the relative pn-pair energy, going out from the photon vertex, coincides in the diagrams under consideration for $E_{\gamma}=45...150$ MeV.

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1. INTRODUCTION

The interest to the problem of three-particle photo disintegration of ${}^{4}He$ nucleus is determined by the possibility of proving the notions about nuclear forces, nucleon-nucleon correlations, cluster configurations and meson exchange currents in nuclei. All these questions are associated with determination of a reaction mechanism. According to the experimental investigations [1] (FIAN, Moscow; Wilson chamber), $\left[2\text{-}4\right]$ (NSC KIPT; diffusion chamber in the magnetic field), [5] (Turino, Italy; diffusion chamber in the magnetic field) and theoretical research [6-12] carried out in 60-70 years using the photon energies E_{γ} =26...150 MeV the main mechanism of ${}^{4}He(\gamma, pn)^{2}H$ reaction is the photon absorption by the pair of nucleons (quasideuteron) being at a short distance from each other. In the 80-th years these investigations were continued at the Kharkov Institute of Physics and Technology using the experimental data obtained by means of a diffusion chamber with an appreciably improved statistics (4.5 thousand of events for the given reaction) [13-15].

A satisfactory fit was obtained for the theoretical estimated values to a pole quasideuteron approximation and the experimental energy distributions of deuterons, protons, neutrons, pair (pn), (pd), (nd) correlations and angular (pn) correlations. In [15] a contribution of triangular ³H, ²H and A=3-Feynman diagrams was analyzed and a little importance of the final state interaction (FSI) even with low energies E_{γ} was stated. However, the interpretation of an effect of the appearance of two maxima at $E_{\gamma} \sim 50$ and 75 MeV in the energy distributions of total crosssections of the $\sigma_t(E_{\gamma})^4 He(\gamma, pn)^2 H$ reaction should be considered as one of unsolved theoretical problems.

Also, unsolved was a problem of two "peaks" appearance in the experimental dN/dE_d energies distri-

butions of deuterons [14] with $E_d^m(I) \approx 1.25$ MeV and $E_d^m(II) \approx 2.5$ MeV at $E_{\gamma}=28...45$ and 28...50 MeV, their contributions to the reaction section being changing with E_{γ} increasing (the first "peak" becomes blurry and the second "peak" becomes distinct). The interpretation of this effect is given in Section 2 within the bounds of extension of the mechanism of ${}^4He(\gamma, pn)^2H$ reaction with taking into account a pole quasi-deuteron mechanism and ${}^3\text{He}({}^3\text{H})$ -triangular diagrams which differ markedly from the triangular diagrams [15] by the (np)-pair knocking-out from the photon vertex.

2. MECHANISMS OF ${}^{4}\text{He}(\gamma, pn)^{2}\text{H}$ REACTION IN THE DIAGRAM REPRESENTATION

The pole mechanism of ${}^{4}He(\gamma, pn)^{2}H$ reaction, being discussed in the introduction, corresponds to the Feynman pole diagram (Fig.1) with a quasi-deuteron photodisintegration outside the mass surface in the photon vertex 1 and with a nuclear form-factor in spectator channel 2.

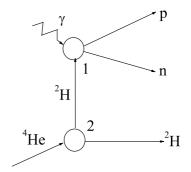


Fig.1. Pole quasideuteron diagram for the ${}^{4}He(\gamma, pn)^{2}H$ reaction

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In the first approximation, to take into account only the pole mechanism features, vertex 1 and vertex 2 of Fig.1 are replaced by the constants. Then the distributions (being the most sensitive to the reaction mechanism) by the relative energy of (pn) pair, $d\Lambda/dt_{pn}$, and the kinetic energy of deuteron, $d\Lambda/dE_d$, take the following form [13,16]

$$\frac{d\Lambda}{dt_{pn}} = C_1 \left[\frac{\varepsilon'}{E_0} + (1 - t_{pn})\right]^{-2} \cdot \frac{dV}{dt_{pn}},\qquad(1)$$

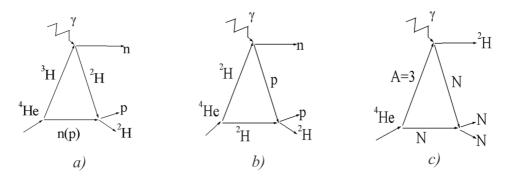
$$\frac{dV}{dt_{pn}} = C_{21} [t_{pn} (1 - t_{pn})]^{1/2}, \qquad (2)$$

$$\frac{d\Lambda}{dE_d} = C_3 [\varepsilon' \frac{m_d}{m_{He}} + E_d]^{-2} \cdot \frac{dV}{dE_d}, \qquad (3)$$

$$\frac{dV}{dE_d} = C_4 [E_d (E_0 - E_d \frac{m_d + m_p + m_n}{m_p + m_n})]^{1/2}, \quad (4)$$

where dV/dt_{pn} and dV/dE_d are the phase distributions, $\varepsilon' = 2m_d - m_{He}$, $E_0 = E_{\gamma} - \varepsilon$, $\varepsilon = m_d + m_p + m_n - m_{He}, t_{pn} = E_{pn}/E_0, E_{pn}$ and E_d are the relative energy of (pn)-pair and the kinetic energy of deuteron in the c.m.s. reaction respectively, $C_1...C_4$ are the constants.

The distributions (1) and (3) correspond to the experimental data [14] and [13] at $E_{\gamma}=30...45$ MeV, 45...75 MeV and $E_{\gamma}=28...150$ MeV respectively with a normalization performed for the total number of events. The taking into account of triangular diagrams for the ${}^{4}He(\gamma, pn)^{2}H$ reaction is considered in [15] and presented in Fig.2. The calculations of $d\Lambda/dE_p$, $d\Lambda/dE_n$, $d\Lambda/dE_d$ in [15] did not allowed one to describe simultaneously, within the theory domain [17], the experimental energy spectra of all the ${}^{4}He(\gamma, pn){}^{2}H$ reaction products in the region of $E_{\gamma}=26...50$ MeV and 50...150 MeV and opened to question the role of FSI in this reaction.



Triangular diagrams [15] for the ${}^{4}He(\gamma, pn)^{2}H$ reaction Fig.2.

the effect of two "peaks" appearance in dN/dE_d for tion cross-section was varying depending on E_{γ} . $E_d^m = 1.25$ and 2.5 MeV at $E_{\gamma} = 28...45$ MeV and

In [14] one discovered, but did not explained, 28...50 MeV the contributions of which to the reac-

Table	1.	

E_{γ}	2830	3035	3540	4050
$E_d^m(\mathbf{I}), \mathrm{MeV}$	$0.66 {\pm} 0.20$	1.25 ± 0.35	$1.81 {\pm} 0.26$	2.39 ± 0.29
$E_d^m(\mathrm{II}), \mathrm{MeV}$		2.22	2.22	2.22

ing into account the pole ${}^{2}H$ -diagram (Fig.1) with the distribution $d\Lambda/dE_d$ in equation (3) and with taking into account the triangular ${}^{3}He({}^{3}H)$ -diagrams (Fig.3) (being qualitative different from the triangular diagram of Fig.2) with a root singularity in the $p+n \rightarrow^2 H$ vertex on the mass shell at a deuteron kinetic energy E_d^* , according to the theory [17], equal to

$$dE_d^m(II) = m_p + m_n - m_{^2H} = 2.22 \, MeV \,. \tag{5}$$

This values of $E_d^*(II)$, owing to proper singularity's triangular diagram independence on E_{γ} [17], determines the position of the stable second "peak" (5)at $E_{\gamma} \geq 30$ MeV. At the same time, the first "peak", unlike the interpretation in [14], moving with E_{γ} changing, is related with the motion of a maximum

In the present paper this effect was obtained tak- of $E_d^*(I)$ in the distribution $d\Lambda/dE_d$ (3), the position of which is determined by the standard procedure of extremum calculation (3)

$$dE_d^m(I) = \frac{3E_0 + 48 - [9E_0^2 + 96E_0 + (48)^2]^{1/2}}{8}.$$
(6)

In Table 1 given are the calculation results for $E_d^m(\mathbf{I})$ by formula (6) and for $E_d^m(\mathbf{II})$ by formula (5) with taken into account the divergances from the average values of E_0 for the intervals of energy $E_{\gamma}=28...30$, 30...35, 35...40 and 40...50 MeV. In whole, a noncontradictory agreement between the model estimates and the experiment takes place.

It is interesting to compare for the diagrams of Fig.1 and Fig.3 the estimated values of t_{pn}^m in the distributions of maxima $d\Lambda/dt_{pn}$ in the different energy intervals E_{γ} , where $t_{pn} = E_{pn}/E_0$. For the pole ²*H*-diagram formula (1) takes place and for the triangular ³*He*(³*H*)-diagrams we use a general equation for the ⁴*He*(γ , pn)²*H* reaction [1,14].

$$\frac{E_{pn}}{E_0} = 1 - \frac{E_d}{E_d^{max}} \,, \tag{C}$$

where

$$E_d^{max} = E_0 \frac{m_p + m_n}{m_d + m_p + m_n}$$

Taking into consideration (5) we obtain

$$t_{pn}^{m} = \frac{E_{pn}^{m}}{E_{0}} = 1 - \frac{2E_{d}^{m}(II)}{E_{0}}$$

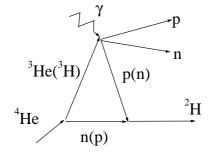


Fig.3. Triangular ${}^{3}He({}^{3}H)$ diagrams for the ${}^{4}He(\gamma, pn){}^{2}H$ reaction

In Table 2 given are the estimated values of $t_{pn}^m(1)$ for the distributions $d\Lambda/dt_{pn}^m$ (1), normalized to all the events, and the estimates of $t_{pn}^m(9)$ by (9) with taking into account the divergences from the average values of E_0 for $E_{\gamma}=30...45$, 45...75 and 75...150 MeV.

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	E_{γ}	3045	4575	75150
ĺ	$t_{pn}^m(1)$	$0.65 {\pm} 0.10$	$0.82 {\pm} 0.05$	$0.92{\pm}0.02$
ĺ	$t_{pn}^m(9)$	$0.61 {\pm} 0.23$	$0.87 {\pm} 0.07$	$0.95 {\pm} 0.03$

As is obvious from Table 2 the estimates of $t_{pn}^m(1)$ and $t_{pn}^m(9)$ are coinciding for $E_{\gamma}=45...150$ MeV. For phase distributions (2) $t_{pn}=0.5$. From the experimental analysis of pn-correlations it is practically impossible to single out the mechanisms of ${}^{4}He(\gamma, pn){}^{2}H$ reactions presented in Figs.1 and 3.

3. DISCUSSION

In the present paper, taking into account the experimental appearance of two "peaks" in the deuteron energy distribution un the range of $E_{\gamma}=28...50$ MeV, for the first time a conclusion (Table 1) was made about the contribution to the ${}^{4}He(\gamma, pn){}^{2}H$ reaction not only from the mechanism of a pole quasideuteron diagram (see Fig.1) with the first "peak"

ole E^m_d(I) moving at an energy E_γ according to Equation (6), but, also, from the mechanism of triangularities and the second string of the second string and the second string second string and the second string second second string second second string second str

(9) Taking into consideration the results obtained upon the mechanisms of ${}^{4}He(\gamma, pn){}^{2}H$ reaction, of particular interest is the experimental and theoretical analysis of residual nuclei (${}^{10}B$, ${}^{12}C$ and ${}^{14}N$) energy distributions in the reactions ${}^{12}C(\gamma, pn){}^{10}B$, ${}^{14}N(\gamma, pn){}^{12}C$ and ${}^{16}O(\gamma, pn){}^{14}N$.

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О МЕХАНИЗМЕ РЕАКЦИИ ${}^{4}He(\gamma, pn)^{2}H$ В.Н. Гурьев

Дана интерпретация эффекта возникновения двух "пиков"в экспериментальных энергетических распределениях дейтронов в реакции ${}^{4}He(\gamma, pn)^{2}H$ при энергиях фотонов $E_{\gamma}=30...50$ МэВ с учетом механизма полюсной ${}^{2}H$ -диаграммы, определяющим движение первого "пика", и механизма треугольных ${}^{3}He({}^{3}H)$ диаграмм с фиксированным проявлением второго пика с энергией дейтрона $E_{2H}^{m}=2.22$ МэВ из-за корневой особенности в вершине $p+n \rightarrow {}^{2}$ Н. Положение максимумов в расчетных распределениях по относительной энергии рп-пары, выходящих из фотонной вершины, совпадает для рассмотренных диаграмм при $E_{\gamma}=45...150$ МэВ.

ПРО МЕХАНІЗМ РЕАКЦІЇ ${}^{4}He(\gamma, pn)^{2}H$ В.М. Гур'ев

Надано інтерпретацію ефекту появи двох "піків"в експериментальних енергетичних розподілах дейтронів в реакції ${}^{4}He(\gamma, pn){}^{2}H$ при енергіях фотонів $E_{\gamma}=30...50$ MeB з урахуванням механізму полюсної ${}^{2}H$ -діаграми, визначаючим рух першого "піку", і механізму трикутних ${}^{3}He({}^{3}H)$ -діаграм з фіксованим проявом другого піку з енергією дейтрона $E_{2H}^{m}=2.22$ MeB через корневу особливість у вершині $p + n \rightarrow {}^{2}H$. Положення максимумів у розрахункових розподілах по відносній енергії рп-пари, що виходять з фотонної вершини, співпадає для розглянутих діаграм при $E_{\gamma}=45...150$ MeB.